Результаты эксперимента LHCb

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LHCb precision measurements



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Outline

Main goal of this talk: Show how precise LHCb measurements in b- and c-sectors make constraints on fundamental prameters of Standard Model (SM) and provide a New Physics (NP) searches.

- Standard Model (SM) and its difficulties
 - Cabibbo-Kobayashi-Maskawa (CKM) matrix, CP violation (CPV)
 - Why and where to find New Physics (NP)? MFV or not?
 - Power of indirect measurements
- LHCb setup (apparatus, physical program etc.)
- Selected results
 - Mixing, CP violation, CKM γ in B systems
 - Mixing and CP violation in charm sector
 - Rare decays $(B \rightarrow 2\mu, B \rightarrow K^*2\mu)$
- Summary and Outlook (what can be achieved after upgrade?)

Introduction

Standard Model

Параметр	Значение	
$\alpha_{\rm s}(M_{\rm Z})$	$0,\!114\pm0,\!0007$	
$1/\alpha(M_{\rm Z})$	$127{,}916\pm0{,}015$	
$\sin^2 \theta_{\mathbf{W}}(M_{\mathbf{Z}})$	$0,23108\pm 0,00005$	
Θ	$ \le 10^{-10} $	
<i>m</i> _u (2 ГэВ)	$2,5^{+0,8}_{-1,0}$ M \ni B	
m _d (2 ГэВ)	$5,0^{+1,0}_{-1,5}$ M $_{2}B$	
m _s (2 ГэВ)	105 ⁺²⁵ ₋₃₅ МэВ	
$m_{\rm c}(m_{\rm c})$	1,266 ^{+0,031} 0,036 ГэВ	
$m_{\rm b}(m_{\rm b})$	$4,\!198\pm0,\!023$ ГэВ	
$m_{\rm t}(m_{\rm t})$	173,10 \pm 1,35 ГэВ	
me	510,998910 ± 0,000013 кэВ	
m_{μ}	$105{,}658367 \pm 0{,}000004{\rm M}{\ni}{\rm B}$	
m_{τ}	1,77682 \pm 0,00016 ГэВ	
θ_{12}	$13,02^\circ\pm0,05^\circ$	
θ_{23}	$2,35^\circ\pm0,06^\circ$	
θ_{13}	$0,199^\circ\pm0,011^\circ$	
δ	$1,20\pm0,08$	
$v(m_{\mu})$	246,221 \pm 0,002 ГэВ	
$M_{ m H}$	115,5–127,0 ГэВ (уровень достоверности 95 %)	

No doubt that SM is great achievement!

(no conflict with HEP)

Reasons for NP:

- 1) Neutrino sector
 - mass
 - oscillations
- 2) Radiative correction to M(Higgs)
 - fine tuning
 - desert between $\boldsymbol{M}_{_{\! EW}}$ and $\boldsymbol{M}_{_{\! GUT}}$
- 3) Astrophysics
 - dark matter
 - baryon asymmetry of Universe

(CPV is needed)

SUSY good candidate to solve 2) & 3)

Cabibbo-Kobayashi-Maskawa

- Flavour eigenstates do not coincide with week eigenstates
- Mixing matrix V_{CKM}
- CP violating phase appears than there are 3 generations of fermions
- Elements of the CKM matrix appear at the decay vertexes

$$\begin{pmatrix} d'\\s'\\b' \end{pmatrix} = \begin{pmatrix} V_{ud} & V_{us} & V_{ub}\\V_{cd} & V_{cs} & V_{cb}\\V_{td} & V_{ts} & V_{tb} \end{pmatrix} \begin{pmatrix} d\\s\\b \end{pmatrix} = V_{CKM} \begin{pmatrix} d\\s\\b \end{pmatrix}$$



Wolfenstein parametrization

$$V_{\text{CKM}} = \begin{pmatrix} 1 - \frac{\lambda^2}{2} & \lambda & A\lambda^3(\varrho - i\eta) \\ -\lambda & 1 - \frac{\lambda^2}{2} & A\lambda^2 \\ A\lambda^3(1 - \varrho - i\eta) & -A\lambda^2 & 1 \end{pmatrix} + O(\lambda^4), \qquad s_{12} = \lambda, \quad s_{23} = A\lambda^2, \quad s_{13} \exp(-i\delta) = A\lambda^3(\rho - i\eta) \\ s_{12} = \lambda = 0,222 \pm 0,002, \quad s_{23} = O(10^{-2}), \quad s_{13} = O(10^{-3}) \end{pmatrix}$$

Unitarity of CKM matrix

Unitarity is a very important

property of
$$V_{CKM}$$

 $v_{CKM}^+ V_{CKM} = V_{CKM} V_{CKM}^+ = 1$ $VV^+ = \sum_l V_{al} V_{bl}^* = \delta_{ab}$,
 $a, b = u, c, t \quad l, m = d, s, b$ $V^+V = \sum_a V_{al}^* V_{am} = \delta_{lm}$,

Neutral Currents can be written in the

first order of EW theory as:

$$\begin{split} &(\bar{\mathbf{u}}\,\bar{\mathbf{c}}\,\bar{\mathbf{t}}\,)\,\gamma_{\mu}(a+c\gamma_{5})\begin{pmatrix}\mathbf{u}\\\mathbf{c}\\\mathbf{t}\end{pmatrix}+(\bar{\mathbf{d}}\,'\bar{\mathbf{s}}\,'\bar{\mathbf{b}}\,')\,\gamma_{\mu}(a+c\gamma_{5})\begin{pmatrix}\mathbf{d}\,'\\\mathbf{s}\,'\\\mathbf{b}\,'\end{pmatrix}=\\ &=(\bar{\mathbf{u}}\,\bar{\mathbf{c}}\,\bar{\mathbf{t}}\,)\,\gamma_{\mu}(a+c\gamma_{5})\begin{pmatrix}\mathbf{u}\\\mathbf{c}\\\mathbf{t}\end{pmatrix}+(\bar{\mathbf{d}}\,\bar{\mathbf{s}}\,\bar{\mathbf{b}})\,\gamma_{\mu}(a+c\gamma_{5})V_{\mathrm{CKM}}^{+}V_{\mathrm{CKM}}\begin{pmatrix}\mathbf{d}\\\mathbf{s}\\\mathbf{b}\end{pmatrix}. \end{split}$$

As a result Flavour Changing Neutral Currents (FCNC) is forbidden at tree level but can appear at the loop level!



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For example :

$\begin{array}{l} & Unitarity \ triangles \\ & V_{ud}V_{cd}^{*} + V_{us}V_{cs}^{*} + V_{ub}V_{cb}^{*} = 0, \\ & V_{ud}V_{td}^{*} + V_{us}V_{ts}^{*} + V_{ub}V_{tb}^{*} = 0, \\ & V_{cd}V_{td}^{*} + V_{cs}V_{ts}^{*} + V_{cb}V_{tb}^{*} = 0, \\ & V_{ud}V_{us}^{*} + V_{cd}V_{cs}^{*} + V_{td}V_{ts}^{*} = 0, \\ & V_{ud}V_{us}^{*} + V_{cd}V_{cs}^{*} + V_{td}V_{ts}^{*} = 0, \\ & V_{ud}V_{ub}^{*} + V_{cd}V_{cb}^{*} + V_{td}V_{tb}^{*} = 0, \\ & V_{us}V_{ub}^{*} + V_{cs}V_{cb}^{*} + V_{ts}V_{tb}^{*} = 0. \end{array}$

- Two of six relation have all three contribution of same size
- Can be drawn as triangle at the complex plane
- Almost all SM CPV is sitting here
- Parameters of the triangle can be measured at the decays





$$\bar{\rho} = \rho\left(1 - \frac{\lambda^2}{2}\right), \quad \bar{\eta} = \eta\left(1 - \frac{\lambda^2}{2}\right)$$

- Many different experimental constraints
- There are another fitting groups
- In this talk I will show LHCb results on γ , Δm





- Other triangles are also very important
- In this talk I will show LHCb results on φs
- Note: Another consequence of V_{CKM} unitarity is that squares of elements in raws and columns must be equal to one. [see for example RPP 73, 046301]

NP and flavour symmetry; Wilson's coefficients

- Progress of theory calculations allows to take into account QCD corrections needed for SM FCNC implementation to decays. (Calculation of C_i in SM as well as quite precise predictions for certain processes)
- H_{eff} is an effective way to test different classes of possible NPs, because C_i depend on their flavour structures.
- Minimal Flavour Violation (MFV) paradigm: NP has same source of FV as SM => real numbers, same CPV effects, relations like:

 $\frac{\mathrm{BR}(B_{\mathrm{s}} \to \mu^{+}\mu^{-})}{\mathrm{BR}(B_{\mathrm{d}} \to \mu^{+}\mu^{-})} = \frac{\tau_{B_{\mathrm{s}}} f_{B_{\mathrm{s}}}^{2} m_{B_{\mathrm{s}}} |V_{ts}|^{2}}{\tau_{B_{\mathrm{d}}} f_{B_{\mathrm{d}}}^{2} m_{B_{\mathrm{d}}} |V_{td}|^{2}}$

 $\Delta F = 1$ operators in the SM and in MFV

$$\mathcal{H}_{ ext{eff}} = -rac{4\ G_F}{\sqrt{2}} rac{e^2}{16\pi^2} rac{V_{tb} V_{ts}^*}{V_{ts}} \sum_i rac{C_i O_i}{i} + ext{h.c.}$$



 If NP contains additional FV sources of *C_i* become complex as well as new CPV effects might appear!

Indirect measurements at LHC

- How NP related to flavour physics?
- Is NP weakly coupled to flavour sector (MFV) or at very high scale?

Important to have a **probes beyond LHC energies** (direct observation)!

• Better to use processes which are either forbidden either highly suppressed in SM

Flavour Changing Neutral Currents (FCNC) can be such a probe



- Other possibilities Lepton Flavour Violation (LFV) m_{LQ}>100TeV [not discussed here, but see LHCb result on B→eµ in PRL 111, 141801]
- CPV in charm sector

Power of indirect measurements

Example #1: CP violation in kaon system

Has been done when only 3 quark were known

1972 Kabayashi-Maskawa 6-quark model

~13 years before Upsilon discovery

Example #2: Weak neutral current (Gargamelle bubble chamber)

 ~ 10 years before Z discovery at UA1/2

Example #3: ARGUS collaboration report large B-mixing

Suggest large mass of top quark

~8 years t has been discovered at Tevatron

LHCb features

Beauty and charm production





• LHCb: forward spectrometer $2 < \eta < 5$

(ATLAS & CMS: $|\eta| \le 2.5$)



• In LHCb acceptance (*pp*-collisions $\sqrt{s} = 7 \text{TeV}$)

 $\sigma(b\bar{b}) = 75.3 \pm 5.4 \pm 13.0 \ \mu b$ Phys.Lett.B694 (2010) 209-216

 $\sigma(c\bar{c}) = 1419 \pm 12 \pm 116 \ \mu b \sim 20 \times \sigma(b\bar{b})$ Largest charm samples in the world Nucl.Phys.B871 (2013) 1

Experimental setup



Operation in 2010/12



- LHCb operates with high efficiency
- Take data at constant instantaneous luminosity rate: $\mathcal{L} \approx 4 \times 10^{32} \, cm^{-2} s^{-1}$

(factor 2 larger than design luminosity)

 Visible pp interactions per bunch crossing μ = 1.7 (50 ns bunch spacing)



We also have set of *pp* data at $\sqrt{s} = 2.76$ TeV (collected in 2011)

LHCb data analysis

Efficient trigger (L0/HLT1/HLT2): 40MHz → 5kHz

Tagging if needed

Event selection

Kinematical and topological info (pT, p, IP, vertex and track quality)

PID information

Cut based or multivariate selection BDT, Neurobayes, *etc*.

Optimization of selection

Using MC Using small sample of real data

Angular analysis++

Check for systematics

And a lot of other checks!



Physics program of LHCb

GOAL: Search for evidence of NP in CP violation and rare decays of beauty and charm hadrons.

(Probing large mass scales via study of virtual quantum loops of new particles)

LHCb results are available in more that 160 papers submitted to journals and 110 conference https://cds.cern.ch/collection/LHCb%20Conference%20Contributions?In=en contributions https://cds.cern.ch/collection/LHCb%20Papers?In=en

Main direction of searches:

- Partially covered in this talk 1) Rare decays **RD** with di-muons 2) Properties of the B systems CPV, Δms ; Γs , $\Delta \Gamma$, ϕs ; CKM y determination 3) Mixing and CPV in the D systems Mixing observ., $\Delta A(CP)$
 - 4) Spectroscopy and production of heavy quarks
 - 5) Electroweak physics
 - 6) Soft QCD physics, pA and Ap results

Not covered =(

Properties of the $B(B^+, B^0, B_s)$ systems

- 1) Direct CP asymmetry in $B_{(s)}^{0}$ decays
- 2) B_{s}^{0} oscillation frequency measurement
- 3) Mixing induced CPV in B_s^0 , e.g. $B_s^0 \rightarrow J/\psi\phi$ and $B_s^0 \rightarrow J/\psi f$ (980)
- 4) CKM angle y



Oscillation frequency for B

Corresponding SM box-diagrams





 $\Delta m_s = 17.768 \pm 0.023 \, (stat) \pm 0.006 \, (syst) \, ps^{-1}$

B: Fast oscillations

Excellent time resolution required!

$$\Gamma = (\Gamma_L + \Gamma_H) / 2;$$

 $\Delta m_s = M_H - M_L$



 $x = (M_H - M_L) / \Gamma;$ $y = (\Gamma_L - \Gamma_H) / 2\Gamma$ Measure time dependent decay rate of

• $PDF \propto \left[e^{-\Gamma t} \cdot \left(\cosh\left(\frac{\Delta\Gamma}{2}t\right) \pm D\cos(\Delta m t) \right) \right] \otimes R(\sigma_t)$ event-by-event

• Mean decay time resolution 44 fs

Most precise measurement up to date Agreement with world average & SM

Also measured in semileptonic decays [arXiv:1308.1302] !

decay time

resolution

Mixing induced CP violation in B_{a}

- Decay of particle and antiparticle to same state
- **CP violating phase** predicted to be **very small in SM** CKMfitter group [PRD 84, 033005]

 $\phi_s^{SM} = -2\beta_s = (-0.0363 \pm 0.0016) \,\mathrm{rad}$

- Observable very sensitive to NP !
- LHCb measured it in two modes (1 fb⁻¹ dataset)
 [PRD 87,112010]
- Measurement of time-dependent CP asymmetry

 $A_{\rm CP}(t) \sim (1 - 2\omega_{\rm tag}) D(\sigma_t) \sin(\Delta m_s(t)) \sin(\phi_s)$

Tagging and high decay time resolution required!





Mixing induced CP violation in B_{s}



Consistent with SM prediction!

$B^0_{s} \rightarrow J/\psi \pi^+ \pi^-$

- Dominated by $f_0 \rightarrow \pi^+ \pi^-$
- BF~35% of $B^{\theta}_{s} \rightarrow J/\psi \phi$
- CP-odd final state
- No angular analysis is required
- Constrain Γ_s and $\Delta \Gamma_s$ to $B^{\theta}_s \to J/\psi \phi$ result

[PLB 713, 378] $\phi_s = -0.019^{+0.173+0.004}_{-0.174-0.003}$ rad.

Combined fit of $B^0 \rightarrow J/\psi \phi$ and $J/\psi \pi^+\pi^-$

 $\phi_s = 0.01 \pm 0.07 (\text{stat}) \pm 0.01 (\text{syst})$ rad,

 $\Gamma_s = 0.661 \pm 0.004 (\text{stat}) \pm 0.006 (\text{syst}) \text{ ps}^{-1}$

 $\Delta\Gamma_s = 0.106 \pm 0.011(\text{stat}) \pm 0.007(\text{syst}) \text{ ps}^{-1}$

[PRD 87, 112010]

Constrain on NP parameters

Consistent with SM prediction

and data from other experiments!



Parameters of CKM triangle

CKM angle y measured with high uncertainty!

$$\gamma = \arg\left[-V_{ud}V_{ub}^*/(V_{cd}V_{cb}^*)\right]$$

(but very precise SM prediction for these observable)



 $\delta \gamma / \gamma < \mathcal{O}(10^{-6})$

GLW / ADS / GGLZ methods

Gronau-London-Wyler (GLW) D in CP-eigenstate $(D \rightarrow KK, \pi\pi)$

[PLB 265, 172 (1991)]

R _{cn} , –	$2[\Gamma(B^- \to D_{CP\pm}K^-) + \Gamma(B^+ \to D_{CP\pm}K^+)]$
$m_{CP\pm} =$	$\Gamma(B^- \to D^0 K^-) + \Gamma(B^+ \to \overline{D}{}^0 K^+)$
Acres -	$\Gamma(B^- \to D_{CP\pm}K^-) - \Gamma(B^+ \to D_{CP\pm}K^+)$
$A_{CP\pm} =$	$\Gamma(B^- \to D_{CP\pm}K^-) + \Gamma(B^+ \to D_{CP\pm}K^+)$.

Atwood-Dunietz-Sony (ADS)

D Cabibbo-allowed $(D^0\to K^-\pi^+)$ and doubly Cabibbo-suppressed $(D^0\to K^+\pi^-)$ states.

$$\begin{split} R_{\rm ADS} &= \frac{\Gamma(B^- \to D[\to \pi^- K^+]K^-) + \Gamma(B^+ \to D[\to \pi^+ K^-]K^+)}{\Gamma(B^- \to D[\to K^- \pi^+]K^-) + \Gamma(B^+ \to D[\to K^+ \pi^-]K^+)} \\ A_{\rm ADS} &= \frac{\Gamma(B^- \to D[\to \pi^- K^+]K^-) - \Gamma(B^+ \to D[\to \pi^+ K^-]K^+)}{\Gamma(B^- \to D[\to \pi^- K^+]K^-) + \Gamma(B^+ \to D[\to \pi^+ K^-]K^+)} \end{split}$$

Giri, Grossman, Soffer and Zupan (GGSZ) deals with self conjugate 3-body final states :

f = D →K_sππ and K_sKK.

Phys.Rev. D68 (2003) 054018

Strong phase varies over the 3-body phase space.

$$x_{\pm} = r_B \cos(\delta_B \pm \gamma) \quad y_{\pm} = r_B \sin(\delta_B \pm \gamma)$$
$$N_{\pm i}^+ = h_{B^+}[K_{\mp i} + (x_+^2 + y_+^2)K_{\pm i} + 2\sqrt{K_i K_{-i}}(x_+ c_{\pm i} \mp y_+ s_{\pm i})]$$
$$N_{\pm i}^- = h_{B^-}[K_{\pm i} + (x_-^2 + y_-^2)K_{\pm i} + 2\sqrt{K_i K_{-i}}(x_- c_{\pm i} \pm y_- s_{\pm i})]$$

$$R_{CP\pm} = 1 + r_B^2 \pm 2r_B \cos \delta_B \cos \gamma$$
$$A_{CP\pm} = \frac{\pm 2r_B \sin \delta_B \sin \gamma}{R_{CP\pm}}.$$

[PRL 78, 3257 (1997)]

$$R_{\text{ADS}} = r_B^2 + r_D^2 + 2r_B r_D \cos\gamma \cos(\delta_B + \delta_D)$$
$$A_{\text{ADS}} = 2r_B r_D \sin\gamma \sin(\delta_B + \delta_D)/R_{\text{ADS}}$$



Result on CKM y



(Two-body GLW/ADS) : B \rightarrow Dh, D \rightarrow hh [Phys. Lett. B712 (2012) 203] (Four-body ADS) : B \rightarrow Dh, D \rightarrow K $\pi\pi\pi$ [LHCb-PAPER-2012-055; arxiv:1303.4646] (GGSZ) : B \rightarrow DK, D \rightarrow K_shh [Phys. Lett. B718 (2012) 43]

The combined results for $B \rightarrow DK$ decays using 1 fb⁻¹(7 TeV) from GLW/ADS/GGSZ plus 2 fb⁻¹ (8 TeV) from GGSZ :



Confidence intervals

$\gamma \in [43.9, 89.5]^{o}$	at	$95\%\mathrm{CL}$
$\gamma \in [55.1, 79.1]^{\rm o}$	\mathbf{at}	$68\%{ m CL}$

Best fit value

$$\gamma = (67 \pm 12)^{\rm o}$$
 at $68\%\,{\rm CL}$



LHCb-CONF-2013-004

LHCb-CONF-2013-006

Mixing and CPV in charm sector

D⁰ mixing



D⁰ mixing



LHCb already reported about first observation of D^{θ} mixing (by single experiment, 9 σ) **PRL 110, 101802**

Newest results on D⁰ mixing (and CPV)

Wrong-sign-to-Right-sign ratio:



Result of the fit with no-CPV assumption:

R_D [10 ⁻³]	$3.568 \pm 0.058 \pm 0.033$
$y' [10^{-3}]$	$4.8 \pm 0.8 \pm 0.5$
x'^2 [10 ⁻⁵]	$5.5 \pm 4.2 \pm 2.6$
χ^2/ndf	86.4/101

 $\mathbf{R}^+ = \mathbf{R}(\mathbf{t})$ WS-to-RS ratio for $\mathbf{D}^0 \rightarrow \mathbf{K}^+ \pi^-$ decay



In case of no-CPV and no-Mixing assumption should be constant!

World-best result!

Newest results on D⁰ mixing (and CPV)

Wrong-sign-to-Right-sign ratio:

$$R(t) \approx R_D + \sqrt{R_D} y' \frac{t}{\tau} + \frac{x'^2 + y'^2}{4} \left(\frac{t}{\tau}\right)^2$$



Result of the fit with no-CPV assumption:

R_D [10 ⁻³]	$3.568 \pm 0.058 \pm 0.033$
$y' [10^{-3}]$	$4.8 \pm 0.8 \pm 0.5$
x'^2 [10 ⁻⁵]	$5.5 \pm 4.2 \pm 2.6$
χ^2/ndf	86.4/101

World-best result!

Fit with CPV assumptions has been also done

Direct CPV: $A_D \equiv (R_D^+ - R_D^-)/(R_D^+ + R_D^-)$ Mixing-Induced $x'^{\pm} = |q/p|^{\pm 1}(x'\cos\phi \pm y'\sin\phi)$ $y'^{\pm} = |q/p|^{\pm 1}(y'\cos\phi \mp x'\sin\phi)$



[arXiv:1309.6534]

Newest results on D⁰ mixing (and CPV)



CP violation in D decays

In SM direct CP violation predicted to be small $\sim 10^{-3}$ - 10^{-4}

Access via asymmetry measurement

$$\begin{split} A_{CP}(f;t) &\equiv \frac{\Gamma(D^{0}(t) \to f) - \Gamma(\bar{D}^{0}(t) \to f)}{\Gamma(D^{0}(t) \to f) + \Gamma(\bar{D}^{0}(t) \to f)} = \underbrace{a_{CP}^{dir}(f)}_{\text{CPV in decay}} + \underbrace{\frac{t}{\tau} a_{CP}^{ind}}_{\text{CPV in mixing + interference}} \end{split}$$

LHCb: Time integrated difference of asymmetries

$$\begin{aligned} \Delta A_{CP} &= A_{CP}(K^+K^-) - A_{CP}(\pi^+\pi^-) \\ &= [a_{CP}^{dir}(K^+K^-) - a_{CP}^{dir}(\pi^+\pi^-)] + \frac{\Delta < t >}{\tau} a_{CP}^{ind} \end{aligned}$$

With 0.6fb^{-1} data sample LHCb found 3.5σ evidence of direct CP violation

$$\Delta(\mathcal{A}^{CP}) = \mathcal{A}^{CP}(D^0 \to K^+ K^-) - \mathcal{A}^{CP}(D^0 \to \pi^+ \pi^-$$

= [-0.82 ± 0.21(stat) ± 0.11(syst)]%

Later some indication came from other experiments

Led to discussion: "Is it sign from NP?"



CP violation in D decays

In **SM direct CP violation** predicted to be **small** $\sim 10^{-3} - 10^{-4}$

Access via asymmetry measurement

$$\begin{split} A_{CP}(f;t) &\equiv \frac{\Gamma(D^{0}(t) \to f) - \Gamma(\bar{D}^{0}(t) \to f)}{\Gamma(D^{0}(t) \to f) + \Gamma(\bar{D}^{0}(t) \to f)} = \underbrace{a_{CP}^{dir}(f)}_{\text{CPV in decay}} + \underbrace{\frac{t}{\tau} a_{CP}^{ind}}_{\text{CPV in mixing + interference}} \end{split}$$

LHCb measured **time integrated difference of asymmetries** $\Delta A_{CP} = A_{CP}(K^+K^-) - A_{CP}(\pi^+\pi^-)$ $= [a_{CP}^{dir}(K^+K^-) - a_{CP}^{dir}(\pi^+\pi^-)] + \frac{\Delta < t >}{\tau} a_{CP}^{ind}$

Two complementary analysis with 1 fb⁻¹ data sample



CP violation in D decays

LHCb results:

[LHCb-CONF-2013-004]

- D* tagged sample (preliminary)
 - $\Delta A_{CP} = (-0.34 \pm 0.15 \, (stat) \pm 0.10 \, (sys)) \ \%$
- μ tagged sample [PLB 723,33] $\Delta A_{CP} = (+0.49 \pm 0.30 (stat) \pm 0.14 (sys)) \%$

Consistent with no CPV hypothesis!





HFAG averages:

$$\begin{aligned} a_{CP}^{ind} &= (-0.010 \pm 0.162) \ \% \\ \Delta a_{CP}^{dir} &= (-0.329 \pm 0.121) \ \% \end{aligned}$$

Note: ΔA_{CP} measurements in $D^+ \to \phi \pi^+$ and $D_s^+ \to K_s^0 \pi^+$ are compatible with 0 arXiv:1303.4906, not discussed here

Rare decays

Rare decays $B_{(s)}^{0} \rightarrow \mu^{+} \mu^{-}$





Helicity suppressed in SM [arXiv 1303.3820]

 $\begin{array}{lll} \mathcal{B}(B_s \to \mu^+ \mu^-) &=& (3.25 \pm 0.17) \times 10^{-9} \\ \mathcal{B}(B^0 \to \mu^+ \mu^-) &=& (1.07 \pm 0.10) \times 10^{-10} \end{array}$

ΔΓ correction [PRD 86, 014027]

$$\begin{split} \mathcal{B}(B_s \to \mu^+ \mu^-)_{} \\ &= \frac{1 + \mathcal{A}^{\mu\mu}_{\Delta\Gamma} . \Delta\Gamma_s / 2\Gamma_s}{1 - (\Delta\Gamma_s / 2\Gamma_s)^2} . \mathcal{B}(B_s \to \mu^+ \mu^-) \\ &= (3.56 \pm 0.18) \times 10^{-9} \end{split}$$

5% precision SM calculations!

Sensitive to new scalar, pseudoscalar, axial-vector particles in loops

In MSSM:

$$\frac{c_{S,P}^{MSSM^2}}{M_A^4} \propto \frac{m_b^2 m_\mu^2 \tan^6 \beta}{M_A^4}$$

Some words about analysis strategy

- Blind analysis of 3fb⁻¹ of data (full 2011-12 sample)
- Robust selection cuts for reduction of combinatorics
- Boosting Decision Tree (BDT) method using 9 topological variables (to avoid correlation with M_{inv})



- BDT trained on signal and bkg MC
- BDT *calibrated* on data using *B*->*h*⁺*h*⁻ as signal and mass sidebands for bkg.
- 8 BDT bins. In each bin, the compatibility of the observed events with <u>bkg only</u> and <u>SM+bkg</u> hypotheses is calculated.



PRL 111, 101805 Result: first evidence of $B_s^0 \rightarrow \mu^+ \mu^-$



Statistical significance 4σ for B⁰_s signal!

Consistent with SM prediction!

Upper limit for $B^{\theta} \rightarrow \mu^{+} \mu^{-}$

	90% C.L.	95% C.L.
Expected bkg	$3.5 imes 10^{-10}$	$4.4 imes 10^{-10}$
Expected bkg + SM	$4.5 imes 10^{-10}$	$5.4 imes 10^{-10}$
Observed	$6.3 imes 10^{-10}$	$7.4 imes 10^{-10}$

Combination of CMS and LHCb results



First evidence of the decay!

LHCb-CONF-2013-012 CMS-PAS-BPH-13-007

Result vs NP



B-hadron \rightarrow Hadron + $\mu^{+}\mu^{-}$, $D \rightarrow \pi \mu^{+}\mu^{-}$



FCNC processes with a lot of observables

Clear experimental signatures with low background

Well developed SM calculations

NP can be found in

- Rates
- Angular distributions
- Asymmetries

As an example zero-crossing point at forward-backward asymmetry for $B^0 \rightarrow K^* \mu^+ \mu^$ is well predicted within SM and has potential for NP searches.



$b \rightarrow x l^+ l^-$ and $c \rightarrow x l^+ l^-$ menu @ LHCb

A lot of channels = a lot of new (Apr-Sep 2013) results

 $b \rightarrow sl^+l^-$

—	Β ⁰ → Κ*μ ⁺μ [−]	JHEP8(2013)131 / 1308.1340	1st multiD angular analysis
_	$B^{0} \rightarrow K \mu^{+} \mu^{-}$	PRL 110, 031801	CP asymmetry
_	$B^+ \rightarrow K^+ \mu^+ \mu^-$	1308.1707 / 1308.1340	ψ(4160) / CP asymmetry
_	$B^{o} \rightarrow \varphi^{*}\mu^{+}\mu^{-}$	arXiv:1305.2168	1st angular analysis
_	$B^{0} \rightarrow K^{*}e^{+}e^{-}$	JHEP 05,(2013)159	1st evidence in low q2
_	$\Lambda_{b} \rightarrow \Lambda \mu^{+} \mu^{-}$	PLB725, 25	baryons, 1st @ LHC

 $\mathbf{C} \rightarrow \boldsymbol{u} \boldsymbol{l}^{+} \boldsymbol{l}^{-}$

 $- D_{(s)}^{+} \to \pi^{+} \mu^{+} \mu^{-} \text{arXiv:} 1304.6365 \qquad \text{fac}$ $D_{(s)}^{+} \to \pi^{-} \mu^{+} \mu^{+}$

factor ~50 improvement in limit



- Using BDT trained on proxy $B \rightarrow K^* J/\psi$
- Background from upper B sideband
- Choice of variables to avoid biases on angles and $q^2 = m^2(\mu\mu)$
- Final selection from BDT decay time, flight direction, trk/vtx quality, p_{τ} , PID
- BR measured relative to $B \rightarrow K^* J/\psi$

Analysis of $B \rightarrow K^* \mu^+ \mu^-$

- Branching fraction measured differential in q² and 3 decay angles
- Limited statistics: $\phi + \pi$ if $\phi < 0$
- Parametric in 4 angular observables
 F_L, A_{FB}, S₃, A₉, from CP asymmetries and averages of decay amplitudes
- Theoretical uncertainties smaller in angular analysis (hadronic form factors)

$$\frac{1}{\mathrm{d}\Gamma/\mathrm{d}q^2} \frac{\mathrm{d}^4\Gamma}{\mathrm{d}q^2\,\mathrm{d}\cos\theta_\ell\,\mathrm{d}\cos\theta_K\,\mathrm{d}\hat{\phi}} \quad \propto$$

First multi-dimensional

angular analysis

$$\begin{bmatrix} F_{\rm L} \cos^2 \theta_K + \frac{3}{4} (1 - F_{\rm L}) (1 - \cos^2 \theta_K) & - \\ F_{\rm L} \cos^2 \theta_K (2 \cos^2 \theta_\ell - 1) & + \\ \frac{1}{4} (1 - F_{\rm L}) (1 - \cos^2 \theta_K) (2 \cos^2 \theta_\ell - 1) & + \\ S_3 (1 - \cos^2 \theta_K) (1 - \cos^2 \theta_\ell) \cos 2\hat{\phi} & + \\ \frac{4}{3} A_{\rm FB} (1 - \cos^2 \theta_K) \cos \theta_\ell & + \\ A_9 (1 - \cos^2 \theta_K) (1 - \cos^2 \theta_\ell) \sin 2\hat{\phi} \end{bmatrix}$$

Analysis of $B \rightarrow K^* \mu^+ \mu^-$



Analysis of $B \rightarrow K^* \mu^+ \mu^-$



Further analysis of $B \rightarrow K^* \mu^+ \mu^-$

$$\frac{1}{\mathrm{d}\Gamma/\mathrm{d}q^2} \frac{\mathrm{d}^4\Gamma}{\mathrm{d}\cos\theta_\ell \,\mathrm{d}\cos\theta_K \,\mathrm{d}\phi \,\mathrm{d}q^2} = \frac{9}{32\pi} \begin{bmatrix} \frac{3}{4}(1-F_\mathrm{L})\sin^2\theta_K + F_\mathrm{L}\cos^2\theta_K + \frac{1}{4}(1-F_\mathrm{L})\sin^2\theta_K \cos 2\theta_\ell \\ -F_\mathrm{L}\cos^2\theta_K \cos 2\theta_\ell + S_3\sin^2\theta_K \sin^2\theta_\ell \cos 2\phi \\ +S_4\sin 2\theta_K \sin 2\theta_\ell \cos\phi + S_5\sin 2\theta_K \sin\theta_\ell \cos\phi \\ +S_6\sin^2\theta_K \cos\theta_\ell + S_7\sin 2\theta_K \sin\theta_\ell \sin\phi \\ +S_8\sin 2\theta_K \sin 2\theta_\ell \sin\phi + S_9\sin^2\theta_K \sin^2\theta_\ell \sin 2\phi \end{bmatrix},$$

Further analysis of $B \rightarrow K^* \mu^+ \mu^-$

$$\frac{1}{\mathrm{d}\Gamma/\mathrm{d}q^2} \frac{\mathrm{d}^4\Gamma}{\mathrm{d}\cos\theta_\ell \,\mathrm{d}\cos\theta_K \,\mathrm{d}\phi \,\mathrm{d}q^2} = \frac{9}{32\pi} \begin{bmatrix} \frac{3}{4}(1-F_\mathrm{L})\sin^2\theta_K + F_\mathrm{L}\cos^2\theta_K + \frac{1}{4}(1-F_\mathrm{L})\sin^2\theta_K \cos 2\theta_\ell \\ -F_\mathrm{L}\cos^2\theta_K \cos 2\theta_\ell + S_3\sin^2\theta_K \sin^2\theta_\ell \cos 2\phi \\ +S_4\sin 2\theta_K \sin 2\theta_\ell \cos\phi + S_5\sin 2\theta_K \sin\theta_\ell \cos\phi \\ +S_6\sin^2\theta_K \cos\theta_\ell + S_7\sin 2\theta_K \sin\theta_\ell \sin\phi \\ +S_8\sin 2\theta_K \sin 2\theta_\ell \sin\phi + S_9\sin^2\theta_K \sin^2\theta_\ell \sin 2\phi \end{bmatrix},$$

 $P'_{6,8}$ predicted to be small!





Further analysis of $B \rightarrow K^* \mu^+ \mu^-$

$$\frac{1}{d\Gamma/dq^{2}} \frac{d^{4}\Gamma}{d\cos\theta_{\ell} d\cos\theta_{K} d\phi dq^{2}} = \frac{9}{32\pi} \left[\frac{3}{4} (1 - F_{L}) \sin^{2}\theta_{K} + F_{L} \cos^{2}\theta_{K} + \frac{1}{4} (1 - F_{L}) \sin^{2}\theta_{K} \cos 2\theta_{\ell} - F_{L} \cos^{2}\theta_{K} \cos 2\theta_{\ell} + S_{3} \sin^{2}\theta_{K} \sin^{2}\theta_{\ell} \cos 2\phi + S_{4} \sin 2\theta_{K} \sin 2\theta_{\ell} \cos \phi + S_{5} \sin 2\theta_{K} \sin 2\theta_{\ell} \sin \phi + S_{6} \sin^{2}\theta_{K} \cos \theta_{\ell} + S_{7} \sin 2\theta_{K} \sin 2\theta_{\ell} \sin \phi + S_{9} \sin^{2}\theta_{\ell} \sin 2\phi \right],$$

$$P'_{6,8} \text{ predicted to be small!} \qquad [arXiv:1308.1707] \qquad 3.7 \sigma 4.30 < q^{2} < 8.68 \text{ GeV}^{2}/c^{4} + 25 \sigma 1.0 < q^{2} < 6.0 \text{ GeV}^{2}/c^{4} + 25 \sigma 1.0 < q^{2} < 6.0 \text{ GeV}^{2}/c^{4} + 25 \sigma 1.0 < q^{2} < 6.0 \text{ GeV}^{2}/c^{4} + 25 \sigma 1.0 < q^{2} < 6.0 \text{ GeV}^{2}/c^{4} + 25 \sigma 1.0 < q^{2} < 6.0 \text{ GeV}^{2}/c^{4} + 25 \sigma 1.0 < q^{2} < 6.0 \text{ GeV}^{2}/c^{4} + 25 \sigma 1.0 < q^{2} < 6.0 \text{ GeV}^{2}/c^{4} + 25 \sigma 1.0 < q^{2} < 6.0 \text{ GeV}^{2}/c^{4} + 25 \sigma 1.0 < q^{2} < 6.0 \text{ GeV}^{2}/c^{4} + 25 \sigma 1.0 < q^{2} < 6.0 \text{ GeV}^{2}/c^{4} + 25 \sigma 1.0 < q^{2} < 6.0 \text{ GeV}^{2}/c^{4} + 25 \sigma 1.0 < q^{2} < 6.0 \text{ GeV}^{2}/c^{4} + 25 \sigma 1.0 < q^{2} < 6.0 \text{ GeV}^{2}/c^{4} + 25 \sigma 1.0 < q^{2} < 6.0 \text{ GeV}^{2}/c^{4} + 25 \sigma 1.0 < q^{2} < 6.0 \text{ GeV}^{2}/c^{4} + 25 \sigma 1.0 < q^{2} < 6.0 \text{ GeV}^{2}/c^{4} + 25 \sigma 1.0 < q^{2} < 6.0 \text{ GeV}^{2}/c^{4} + 25 \sigma 1.0 < q^{2} < 6.0 \text{ GeV}^{2}/c^{4} + 25 \sigma 1.0 < q^{2} < 6.0 \text{ GeV}^{2}/c^{4} + 25 \sigma 1.0 < q^{2} < 6.0 \text{ GeV}^{2}/c^{4} + 25 \sigma 1.0 < q^{2} < 6.0 \text{ GeV}^{2}/c^{4} + 25 \sigma 1.0 < q^{2} = 6.0 \text{ GeV}^{2}/c^{4} + 25 \sigma 1.0 < q^{2} = 6.0 \text{ GeV}^{2}/c^{4} + 25 \sigma 1.0 < q^{2} = 6.0 \text{ GeV}^{2}/c^{4} + 25 \sigma 1.0 < 10 +$$

Summary

LHCb, the forward spectrometer for precision studies in flavour physics domain

Excellent performance of the LHC and LHCb has led to a lot of physics results

- Test of SM
- Search for NP
- Make CP violation measurements in b- and c-sectors

World best quality of the results in charm and beauty physics!

Remember, that presented here measurements use mainly the 1 fb⁻¹ dataset

(70% of the 2010-12 data still in progress)

OUTLOOK:

1) Plan to have more than ~ 5 fb⁻¹ at $\sqrt{s} = 13$ TeV during next LHC run (2015-18)

= ~ 8 times higher statistics in 2019 (in comparison with presented results)

2) **Upgrade** (next slide)

Outlook. *Theory vs. 50 fb⁻¹*

Туре	Observable	LHCb 2018	Upgrade (50 fb ⁻¹)	Theory uncertainty
B_s^0 mixing	$2\beta_s(B_s^0 \to J/\psi\phi)$	0.025	0.008	~0.003
	$2\beta_s(B_s^0 \rightarrow J/\psi f_0(980))$	0.045	0.014	~ 0.01
	$a_{ m sl}^s$	0.6×10^{-3}	0.2×10^{-3}	0.03×10^{-3}
Gluonic penguins	$2\beta_s^{\rm eff}(B_s^0 \to \phi\phi)$	0.17	0.03	0.02
	$2\beta_s^{\rm eff}(B_s^0 \to K^{*0}\overline{K}^{*0})$	0.13	0.02	< 0.02
	$2\beta^{\rm eff}(B^0 \to \phi K^0_S)$	0.30	0.05	0.02
Right-handed currents	$2\beta_s^{\text{eff}}(B_s^0 \to \phi \gamma)$	0.09	0.02	< 0.01
	$\tau^{\rm eff}(B^0_s\to\phi\gamma)/\tau_{B^0_s}$	5 %	1 %	0.2 %
Electroweak penguins	$S_3(B^0 \to K^{*0} \mu^+ \mu^-; 1 < q^2 < 6 \text{ GeV}^2/c^4)$	0.025	0.008	0.02
	$s_0 A_{\rm FB} (B^0 \rightarrow K^{*0} \mu^+ \mu^-)$	6 %	2 %	7 %
	$A_{\rm I}(K\mu^+\mu^-; 1 < q^2 < 6 { m GeV}^2/c^4)$	0.08	0.025	~ 0.02
	$\mathcal{B}(B^+ \to \pi^+ \mu^+ \mu^-)/\mathcal{B}(B^+ \to K^+ \mu^+ \mu^-)$	8 %	2.5 %	$\sim \! 10 \%$
Higgs penguins	$\mathcal{B}(B^0_s \to \mu^+ \mu^-)$	$0.5 imes 10^{-9}$	$0.15 imes 10^{-9}$	$0.3 imes 10^{-9}$
	$\mathcal{B}(B^0\to\mu^+\mu^-)/\mathcal{B}(B^0_s\to\mu^+\mu^-)$	$\sim 100 \%$	\sim 35 %	$\sim 5~\%$
Unitarity triangle angles	$\gamma(B\to D^{(*)}K^{(*)})$	4°	0.9°	negligible
	$\gamma(B_s^0 \to D_s K)$	11°	2.0°	negligible
	$\beta(B^0 \to J/\psi K_{\rm S}^0)$	0.6°	0.2°	negligible
Charm CP violation	A_{Γ}	0.40×10^{-3}	$0.07 imes 10^{-3}$	-
	ΔA_{CP}	0.65×10^{-3}	$0.12 imes 10^{-3}$	-

Thank you for your attention!

Spare slides

LHCb trigger



Goal: To select interesting beauty and charm decays while maintaining the managable data rates

Level 0:

- Largest $p_{T}(E)$ used
- Typical thresholds 1.5 3.5 GeV/c

Software HLT1:

- Partial event reconstruction
- Selection based on p_{T} , IP

Software HLT2:

- Full event reconstruction
- Mass cuts

On-line charm and strange signals Data quality from sig-to-bkg ratio.



Run 90645, started 2011-05-02 08:06:31, duration: 01:00:24



X(3872) quantum numbers

It is extremely narrow. Only upper limits on its width (<1.2 MeV) None of the known cc states above DD threshold is so narrow

- This automatically eliminates all cc excitations which can decay to DD
- Its mass is not near any of the predicted cc masses. Closest predicted cc states which could be narrow: 2³P₁++, 1¹D₂-+
- Its mass is nearly equal $m(D^0)+m(D^{0*})$:
 - It is loosely bound D⁰D⁰*=(cu)(cu) molecule or (ccuu) tetraquark? Both models require J^{PC}=1⁺⁺

J^{PC}	decay	LS	χ^2 (11 d.o.f.)	χ^2 prob.
1++	$J/\psi \rho^0$	01	13.2	0.28
2^{-+}	$J/\psi \rho^0$	11,12	13.6	0.26
1	$J/\psi(\pi\pi)_S$	01	35.1	2.4×10^{-4}
2+-	$1/\psi(\pi\pi)_s$	11	38.9	5.5×10-5
1+-	$J/\psi(\pi\pi)_S$	11	39.8	3.8×10^{-5}
2	$J/\psi(\pi\pi)_S$	21	39.8	3.8 ×I10-5
3+-	$J/\psi(\pi\pi)_S$	31	39.8	3.8×10^{-5}
3	$J/\psi(\pi\pi)_S$	21	41.0	2.4×10^{-3}
2++	$J/\psi \rho^0$	02	43.0	1.1×10^{-5}
1^{-+}	$J/\psi \rho^0$	10,11,12	45.4	4.1×10^{-6}
0^{-+}	J/400	11	104	3.5×10^{-1}
0+-	$J/\psi(\pi\pi)_S$	11	129	$\leq 1 \times 10^{-2}$
0++	$J/\psi\rho^0$	00	163	$\leq 1 \times 10^{-3}$

CDF's binned 3D angular χ^2 fit:

Cannot distinguish between 1⁺⁺ and 2⁻⁺ All other ruled out.

Discovered by Belle in 2003 at e+e-









- The Gaussian approximation conservative since the actual distribution to the left of the Gaussian fit.
 - The 2⁻⁺ hypothesis is ruled out at 8.4 σ (>8 after systematics)
- 1⁺⁺ C.L. is high (34%).

The state $\eta_{c2}(1^1D_2)$ is excluded, favour unconventional interpretations $\chi_{c1}(2^3P_1)$, $D^{*0}\overline{D}^0$ molecule, tetra quarks or charmonium-molecules

D meson mass measurement

Interpreting X(3872) as $D^{*0}D^0$ molecule E_B is determined by D mass measurements: $E_B = 0.16 \pm 0.26 M eV/c^2$

- Mass measurements in the D system arXiv: 1304.6865 ($\int \mathcal{L} = 1 f b^{-1}$)
 - Determine D^0 mass in $D^0 \to K^+ K^- K^- \pi^+$ $M(D^0) = 1864.75 \pm 0.15 (\text{stat}) \pm 0.11 (\text{sys}) \text{ MeV/c}^2$
 - Mass difference measurements $M(D^+) - M(D^0) = 4.76 \pm 0.12(\text{stat}) \pm 0.07(\text{sys}) \text{ MeV/c}^2$ $M(D_s^+) - M(D^+) = 98.68 \pm 0.03(\text{stat}) \pm 0.04(\text{sys}) \text{ MeV/c}^2$ Derive a significantly more precise D_s^+ mass $M(D_s^+) = 19684.19 \pm 0.20 \pm 0.14 \pm 0.08 MeV/c^2$ CLE02
 - Dominant syst. uncertainty on the mass is due to the momentum scale of 0.03 % D^0 mass $: 0.09 \,\mathrm{MeV/c^2}$ mass difference $: 0.04 \,\mathrm{MeV/c^2}$



Flavour symmetry of SM

 $\mathcal{L}_{\text{SM}} = \mathcal{L}_{\text{gauge}} + \mathcal{L}_{\text{Higgs}} + \mathcal{L}_{\text{Yukawa}}$

- \mathcal{L}_{gauge} and \mathcal{L}_{Higgs} are flavour invariant
- Only \mathcal{L}_{Yukawa} distinguishes flavour (=breaks the flavour symmetry)
- TeV-scale NP as suggested by the hierarchy problem is incompatible with *generic* flavour-violation. This is the **flavour problem** of NP.
- But the size of the *c_i* depends on the *flavour structure* of the NP theory.
- The NP sector has to be approximately invariant under some global flavour symmetry.

The MFV assumption

The SM Yukawas are the only sources of breaking of $U(3)^3$ even beyond the SM

Consequences of MFV

All FCNC amplitudes are governed by the **same** CKM factors as in the SM

- NP effects in b → s, b → d and s → d transitions are perfectly correlated
- Ratios such as

$$\frac{\mathsf{BR}(B_{\mathsf{s}} \to \mu^+ \mu^-)}{\mathsf{BR}(B_{\mathsf{d}} \to \mu^+ \mu^-)} = \frac{\tau_{B_{\mathsf{s}}} f_{B_{\mathsf{s}}}^2 m_{B_{\mathsf{s}}} |V_{ts}|^2}{\tau_{B_{\mathsf{d}}} f_{B_{\mathsf{d}}}^2 m_{B_{\mathsf{d}}} |V_{td}|^2}$$

are not modified and consitute a test of the MFV paradigm

• CP violating phase aligned with the SM \Rightarrow CP asymmetries mostly SM-like

Effective Hamiltonian and Wilson's coefficients

 $\Delta F = 1$ operators in the SM and in MFV

$$\mathcal{H}_{\mathrm{eff}} = -rac{4~G_F}{\sqrt{2}} rac{e^2}{16\pi^2} rac{V_{tb}V_{ts}^*}{V_{ts}} \sum_i rac{C_iO_i}{i} + \mathrm{h.c.}$$

Current-current operators

 $O_1 = (\bar{s}\gamma^{\mu}T^aP_Lc) \otimes (\bar{c}\gamma_{\mu}T^aP_Lb) \qquad O_2 = (\bar{s}\gamma^{\mu}P_Lc) \otimes (\bar{c}\gamma_{\mu}P_Lb)$

QCD penguin operators

$$\begin{split} O_{3} &= \left(\bar{s}\gamma^{\mu}P_{L}b\right) \otimes \sum_{q} \left(\bar{q}\gamma_{\mu}q\right) & O_{4} &= \left(\bar{s}\gamma^{\mu}T^{a}P_{L}b\right) \otimes \sum_{q} \left(\bar{q}\gamma_{\mu}T^{a}q\right) \\ O_{5} &= \left(\bar{s}\gamma^{\mu}\gamma^{\nu}\gamma^{\rho}P_{L}b\right) \otimes \sum_{q} \left(\bar{q}\gamma_{\mu}\gamma_{\nu}\gamma_{\rho}q\right) & O_{6} &= \left(\bar{s}\gamma^{\mu}\gamma^{\nu}\gamma^{\rho}T^{a}P_{L}b\right) \otimes \sum_{q} \left(\bar{q}\gamma_{\mu}\gamma_{\nu}\gamma_{\rho}T^{a}\right) \end{split}$$

Scalar operators

$$O_{S} = \frac{m_{b}}{m_{B_{s}}}(\bar{s}_{L}b_{R})(\bar{\ell}\ell) \qquad O_{P} = \frac{m_{b}}{m_{B_{s}}}(\bar{s}_{L}b_{R})(\bar{\ell}\gamma_{5}\ell)$$

Dipole operators

$$O_7 = \frac{m_b}{e} (\bar{s}_L \sigma_{\mu\nu} b_R) F^{\mu\nu} \qquad O_8 = \frac{g_s m_b}{e^2} (\bar{s}_L \sigma_{\mu\nu} T^a b_R) G^{\mu\nu a}$$

Semileptonic operators

$$O_{9} = (\bar{s}_{L}\gamma_{\mu}b_{L})(\bar{\ell}\gamma^{\mu}\ell) \qquad O_{10} = (\bar{s}_{L}\gamma_{\mu}b_{L})(\bar{\ell}\gamma^{\mu}\gamma_{5}\ell) O_{\nu} = (\bar{s}_{L}\gamma_{\mu}b_{L})(\bar{\nu}_{L}\gamma^{\mu}\nu_{L})$$

NB: in MFV, all Wilson coefficients *C_i* are **real**, dimensionless numbers

 C_{1-6} dominated by QCD contributions and hardly sensitive to NP! $C_{7-10,\nu}$ are generated in the SM and **sensitive** to many **NP** models $C_{S,P}$ **vanish** in the SM but can be sizable in models with extended Higgs sector

Effective Hamiltonian and Wilson coefficients

 $\Delta F = 1$ operators in the SM and in MFV

$$\mathcal{H}_{ ext{eff}} = -rac{4\ G_F}{\sqrt{2}} rac{e^2}{16\pi^2} rac{V_{tb}V_{ts}^*}{V_{ts}} \sum_i rac{C_iO_i}{i} + ext{h.c.}$$

 $b \rightarrow s$ $B \rightarrow X_s \gamma$ 07,8 Y $B \to K^* \gamma$ $B \rightarrow K \ell^+ \ell^ \ell^+\ell^- \quad B \to K^*\ell^+\ell^ O_{7,8}, O_9, O_{10}$ $B \rightarrow X_s \ell^+ \ell^ B_s \rightarrow \mu^+ \mu^ O_{10}, O_{S,P}$ $B \rightarrow X_s \nu \bar{\nu}$ $\nu \overline{\nu}$ $B \rightarrow K \nu \bar{\nu}$ O_{ν} $B \rightarrow K^* \nu \bar{\nu}$

 $B \rightarrow X_{s}\mu\mu$, $B \rightarrow K^{*}\mu\mu$ and $B_{s} \rightarrow \mu\mu$



Beyond MFV

 $\Delta F = 1$ operators in the SM and in MFV

$$\mathcal{H}_{ ext{eff}} = -rac{4~G_F}{\sqrt{2}} rac{e^2}{16\pi^2} rac{V_{tb}V_{ts}^*}{V_{ts}} \sum_i rac{C_iO_i}{i} + ext{h.c.}$$

In $U(2)^3$, no new operators beyond MFV are generated, but Wilson coefficients are now **complex numbers**!

- Correlation between *B* and *K* decays **broken**
- New CPV phase in B decays

 $B \rightarrow X_s \gamma, B \rightarrow (K, K^*, X_s) \mu \mu$ and $B_s \rightarrow \mu \mu$



LHCb result on $B \rightarrow K^* \mu^{\dagger} \mu^{-}$ pull "theory trigger"

- Of course this is only evidence, which must be supported by further experimental analysis
- But everybody can start to think the pointer to which direction of NP it is.
- Here is an example of such studies arXiv:1307.5683 [this is not LHCb result]
- Combination with another world data gives 4.5σ deviation respect to SM
- Possible large NP in C₉.

