

proton-nucleus interactions

V.A. Khoze, M.G. Ryskin – 2307.08625

Advantage of the LRG events

– the background from underlying/MPI events is practically absent.

LRG events are attractive for:

- i) searching the new/BSM physics
- ii) studying the Pomeron-Pomeron (multi-Pomeron) interactions.

1. Kinematics

2. LRG in pp-collisions

a) SD, DD

b) multi-Pomeron exchange/vertices

3. LRG in proton-lead collisions

a) incoherent

gap survival factor

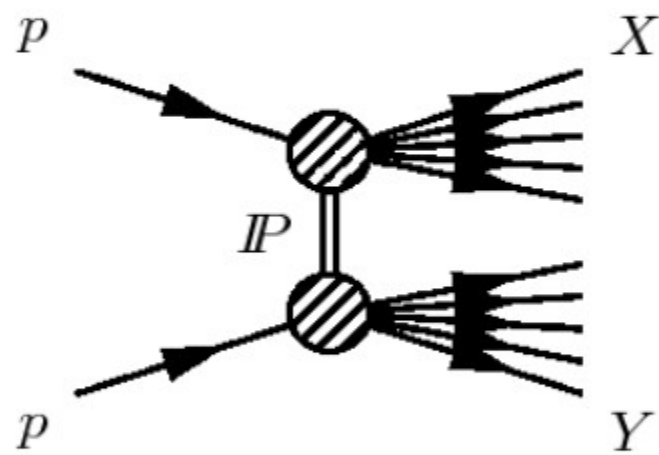
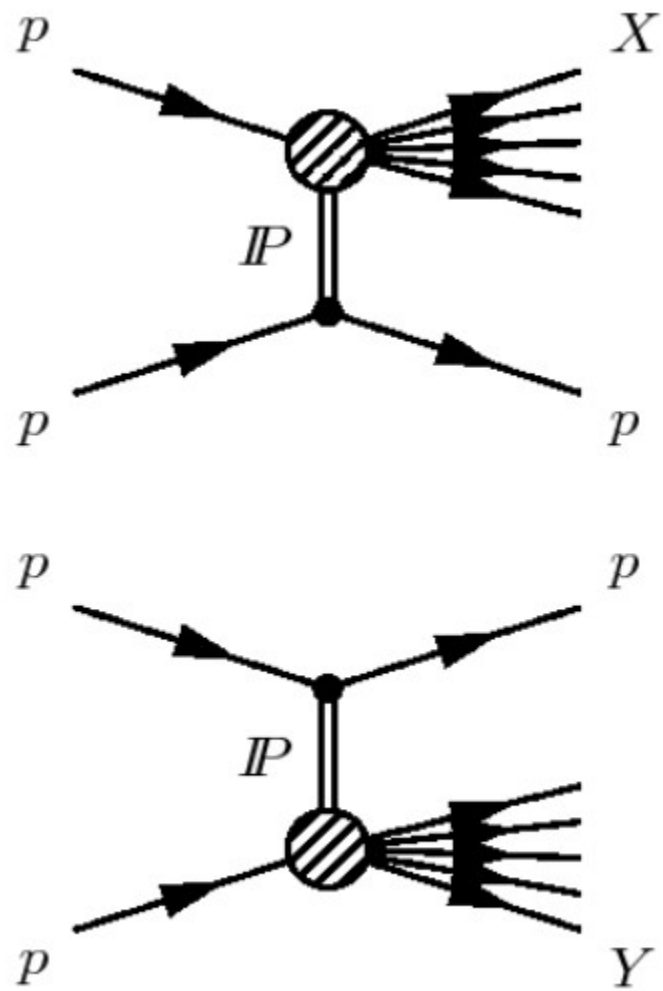
peripheral interaction

role of the LRG pp-amplitude size

b) coherent Pomeron emission

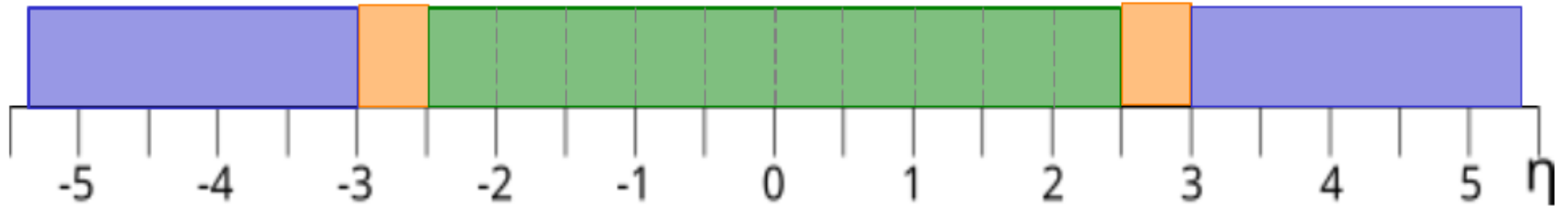
c) photon exchange

4. comparison with the CMS 2301.07630 results.



(c)

Pb-p diffraction [arXiv:2301.07630]: Used thresholds

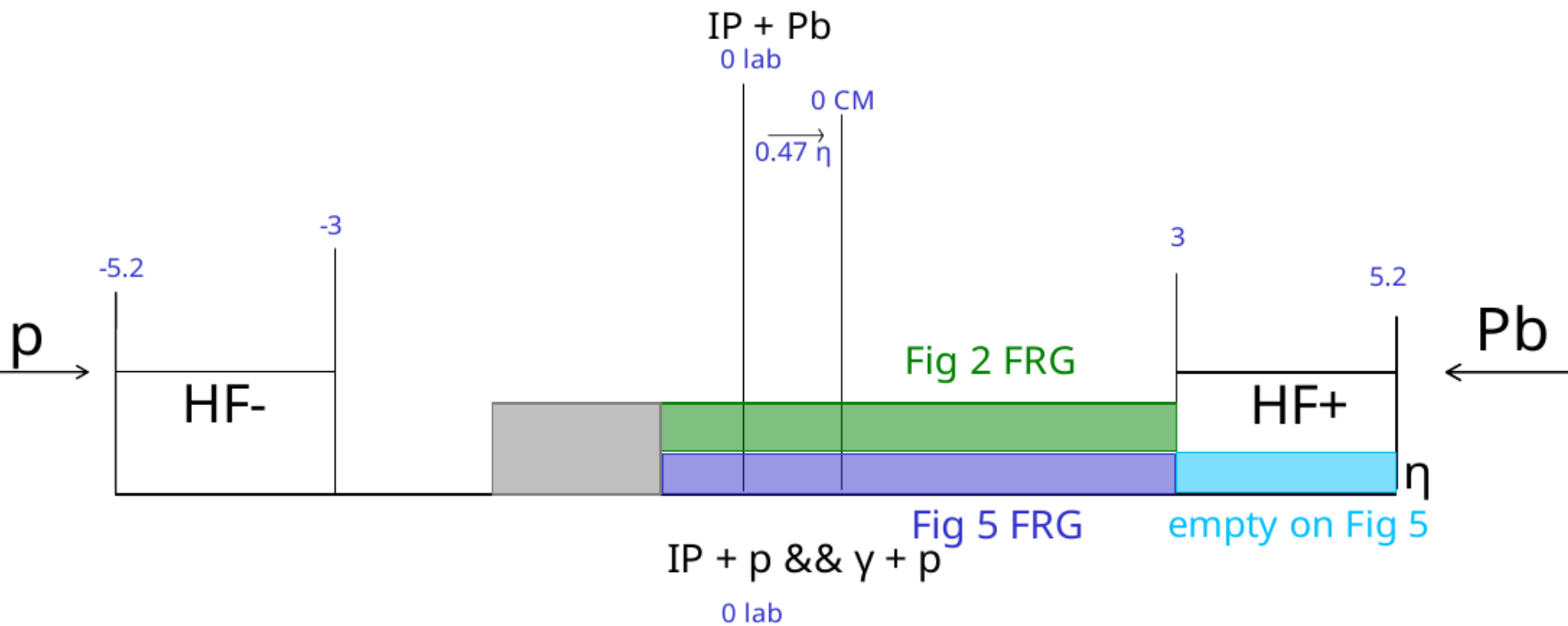
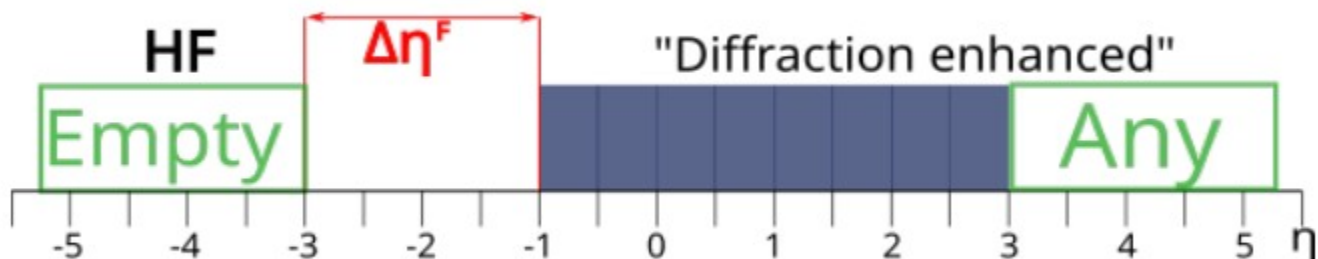
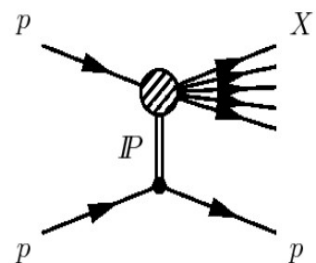
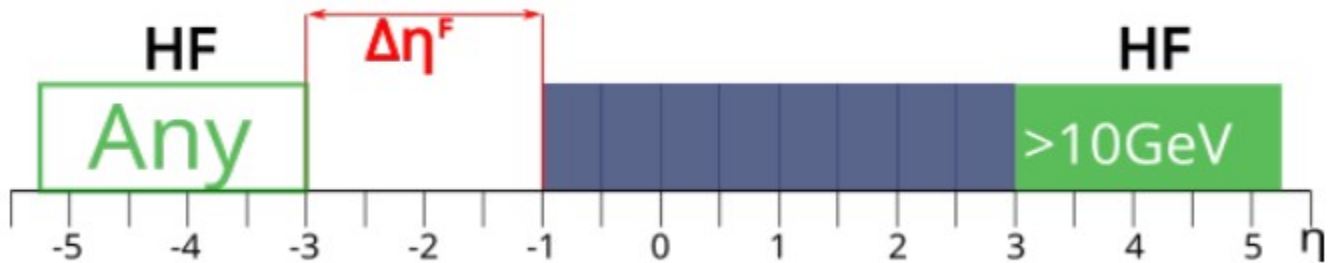
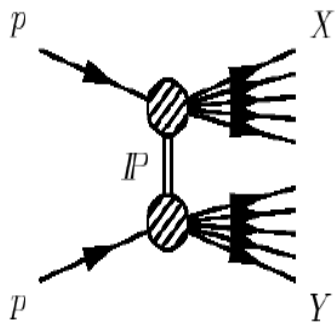


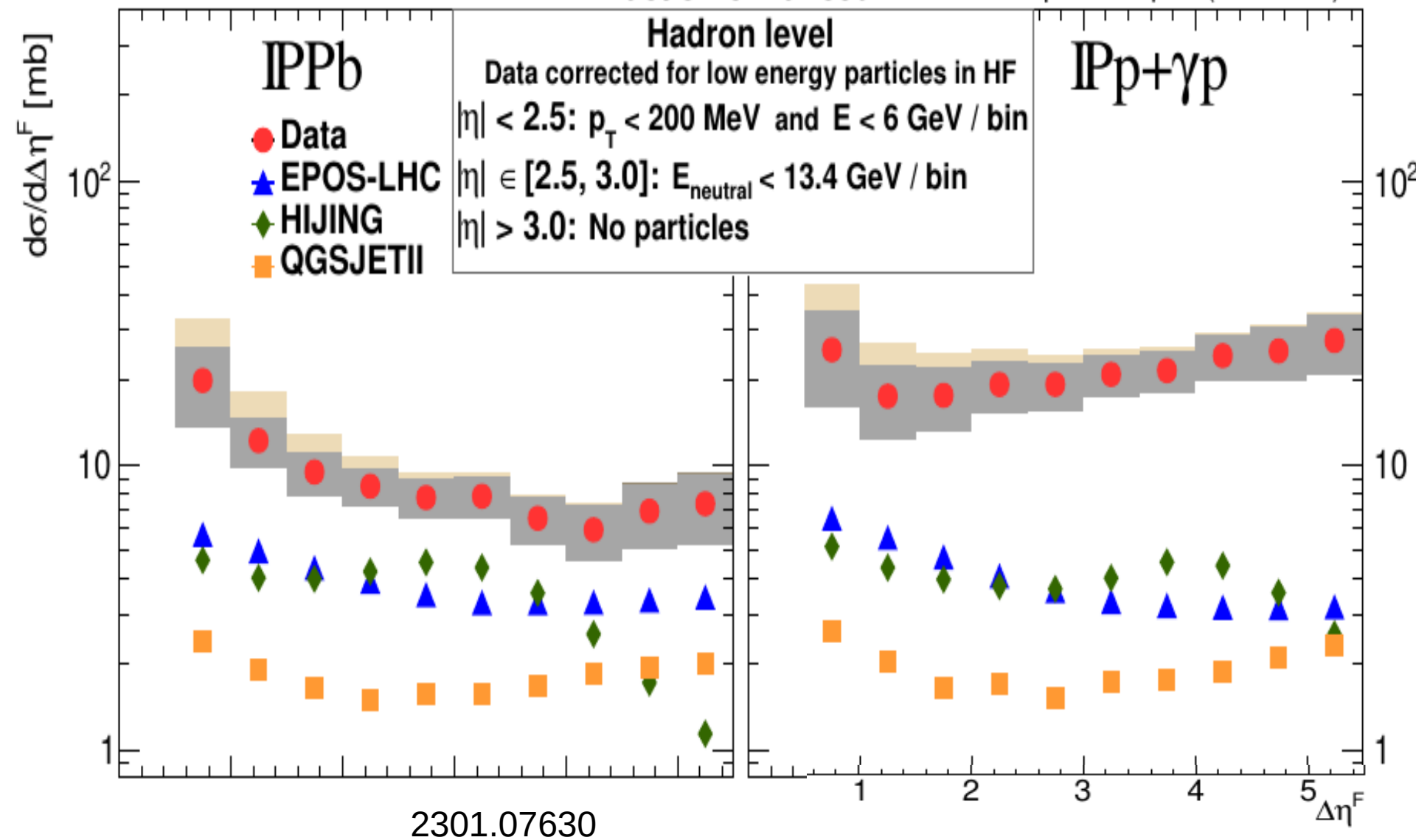
Central part with tracking ($\eta \in [-2.5; 2.5]$) η bin is empty:

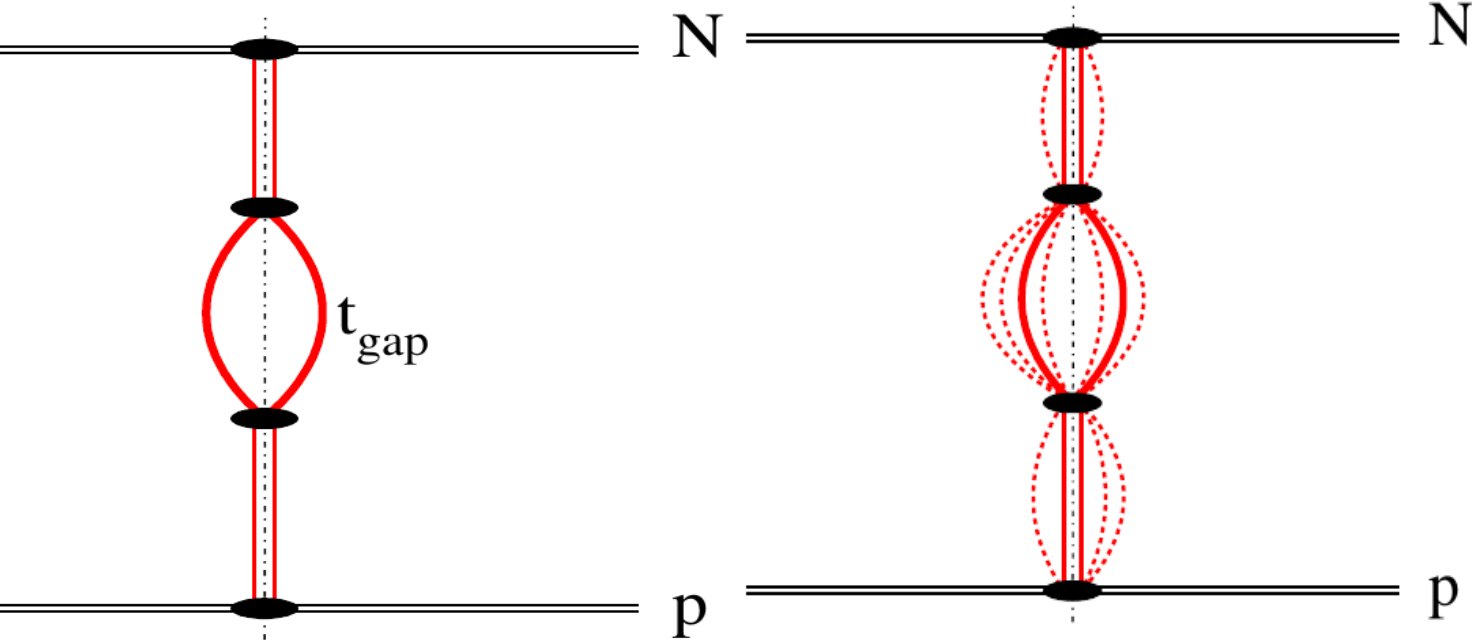
$$\text{no tracks with } p_T > 200 \text{ MeV and } \sum_{bin} E_{pf} < 6 \text{ GeV}$$

Central part without tracking ($|\eta| \in [2.5; 3.0]$) η bin is empty: $\sum_{bin} E_{pf}^{neutral} < 13.4 \text{ GeV}$

HF ($|\eta| \in [3.0; 5.2]$) on the side with RG:



CMS**Diffraction-enhanced**pPb 6.4 μb^{-1} (8.16 TeV)



Pomeron $\Rightarrow \Omega(b, s) = \text{opacity}$

$$\Omega(b, s) = \frac{g_N^2}{2\pi B} \left(\frac{s}{s_0} \right)^{\alpha_P(0)-1} e^{-b^2/2B}$$

$$2\text{Im}A_{el}(b) = |A_{el}(b)|^2 + G_{inel}(b)$$

$$A_{el} = i(1 - e^{-\Omega(b)/2}) \quad G_{inel} = 1 - e^{-\Omega(b)}$$

probability Not to have inelastic interaction is $S^2(b) = e^{-\Omega(b)}$

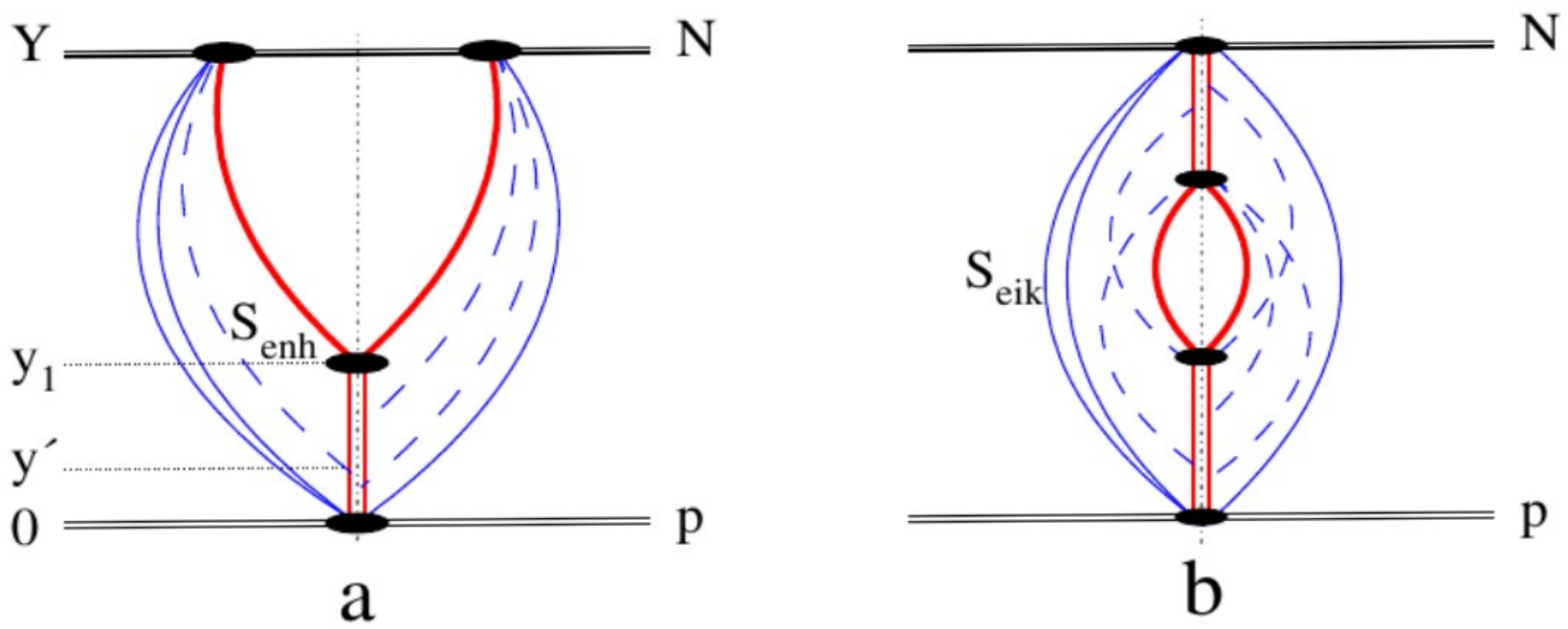


Figure 3: Diagrams for the single (a) and double (b) diffractive dissociation. Heavy red lines denote the elastic amplitude $1 - e^{-\Omega(\Delta b)/2}$ originated by the Pomeron(s) exchange across the LRG; double red lines cut by the dot-dashed line denote the inelastic Pomeron-induced amplitudes $1 - e^{-\Omega(b)}$. In both cases, $\Omega(b)$ is the corresponding opacity caused by one Pomeron exchange in the b representation. Dashed lines describe the semi-enhanced gap survival factor. Here we account for the whole inelastic interactions. That is, the dashed lines correspond to $(1 - e^{-\Omega})$ and integration over the interactions with any intermediate Parton in central pomeron (from $y' = 0$ to $y' = y_1$ in the case of the left figure and analogous for the upper part of the right figure) is included. The continuous thin blue lines indicate the non-enhanced Pomeron exchange. Its sum forms the eikonal survival factor $S_{eik}^2 = \exp(-\Omega(b))$.

$$S_{enh}^2(\delta b) = \exp\left(-\lambda_1 \int_0^{y_1} dy' (1 - e^{\Omega_{parton}(\delta b)})\right)$$

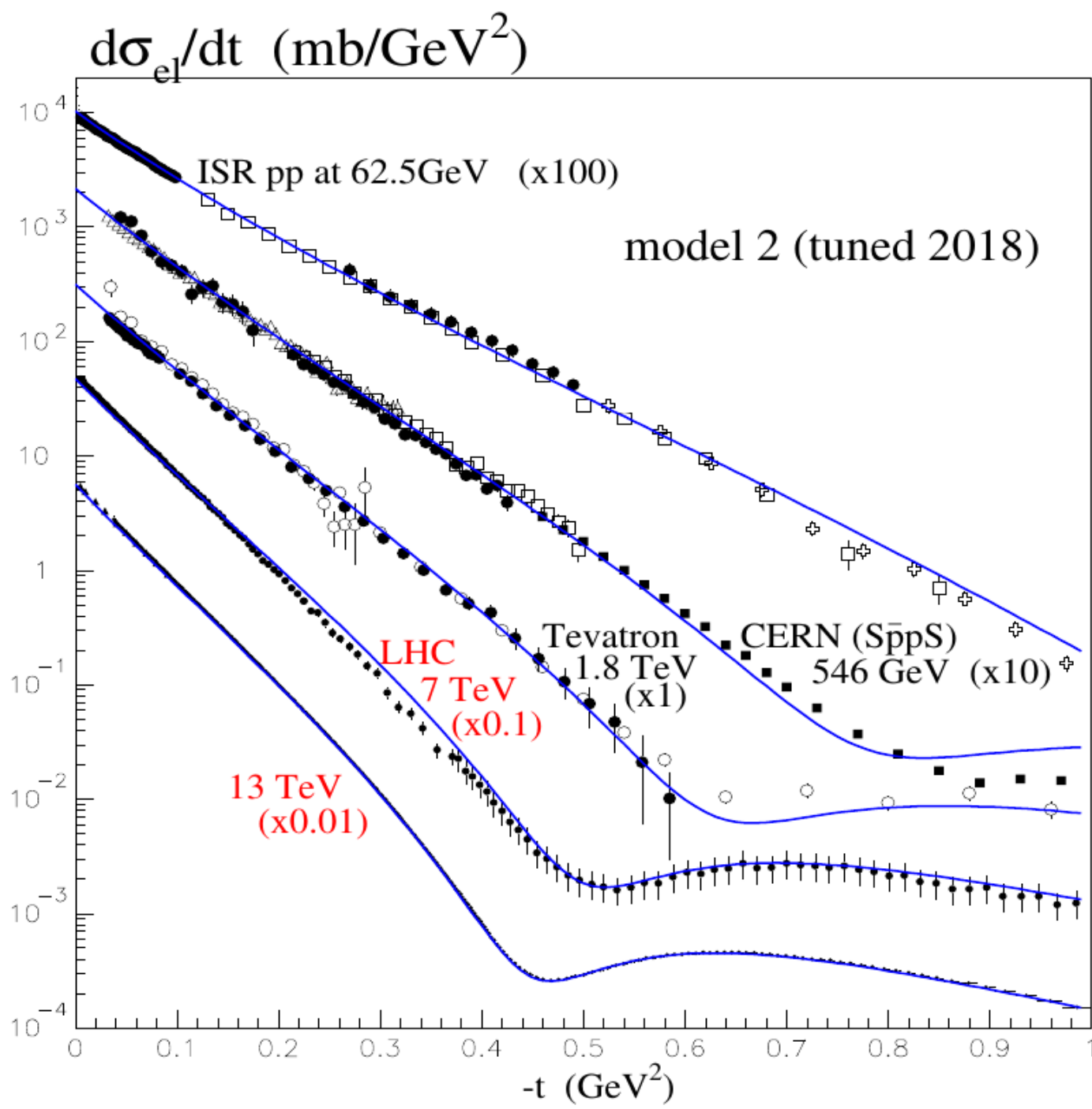
$$g_n^m(k_i^2, k_j^2) = \lambda_1(\lambda_2 g_N)^{m+n-2} \prod_{i=1}^m e^{-b_{3P} k_i^2} \prod_{j=1}^n e^{-b_{3P} k_j^2}, \quad (9)$$

where g_N is the Pomeron nucleon coupling taking from [11], $k_i^2, k_j^2 > 0$ are the transverse momenta squared transferred through the corresponding Pomeron, and we put the slope $b_{3P} = 0.5 \text{ GeV}^{-2}$. $\lambda_1 = 0.4$ and $\lambda_2 = 0.5$, that is $g_{3P} = \lambda_1 \lambda_2 g_N = 0.2 g_N$.

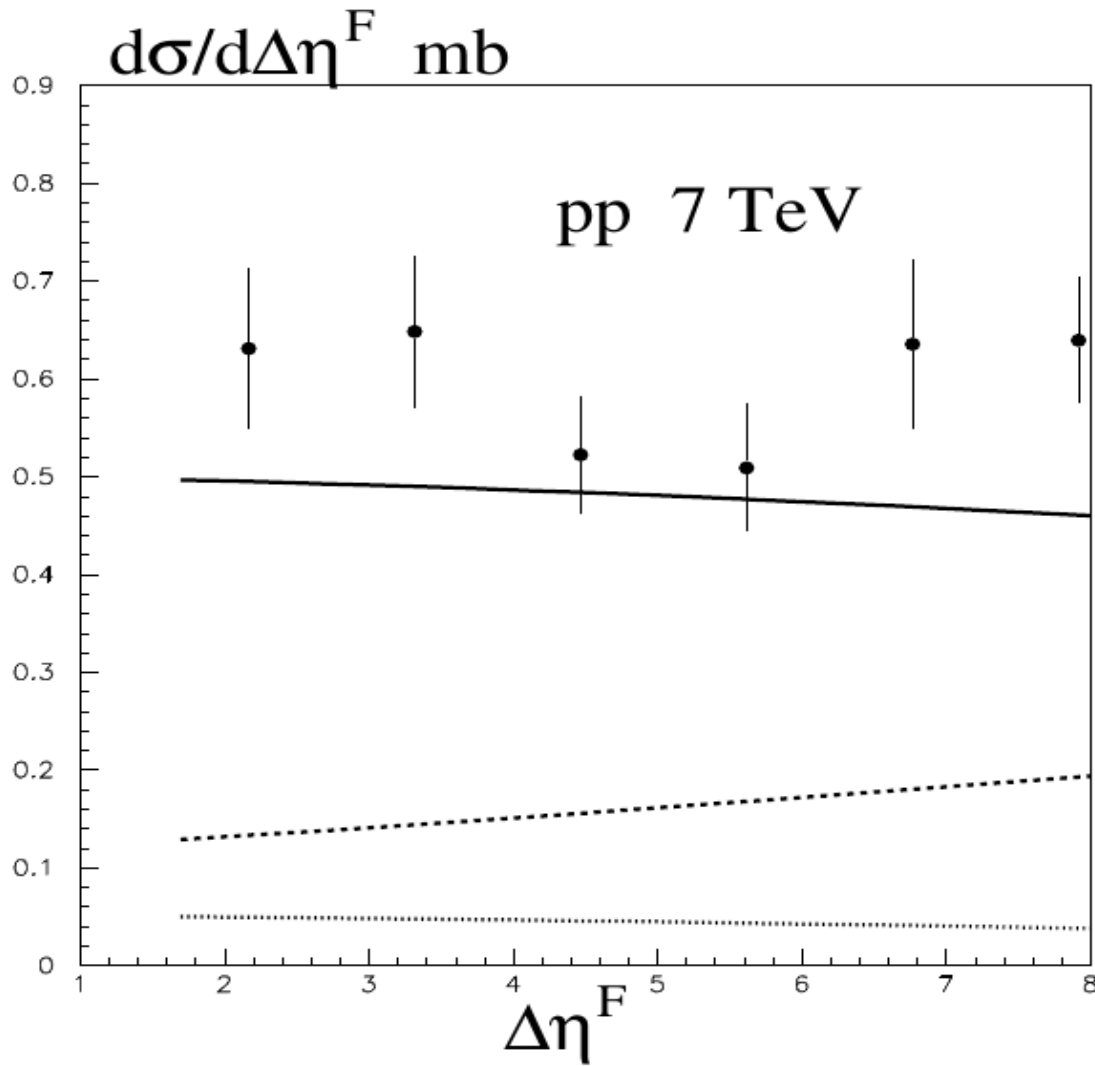
The constant $\lambda_1 = dn/dy$ is the density of the intermediate partons, dn/dy , in the rapidity space while λ_2 is the ratio $g_{parton}/g_N = \lambda_2$ of the Pomeron coupling to the parton at the edge of LRG, g_{parton} , to the Pomeron-nucleon coupling, g_N .

Pomeron parameters are from two (Good-Walker) channel KMR model 1806.05970 $\alpha_P(0) = 1.13$

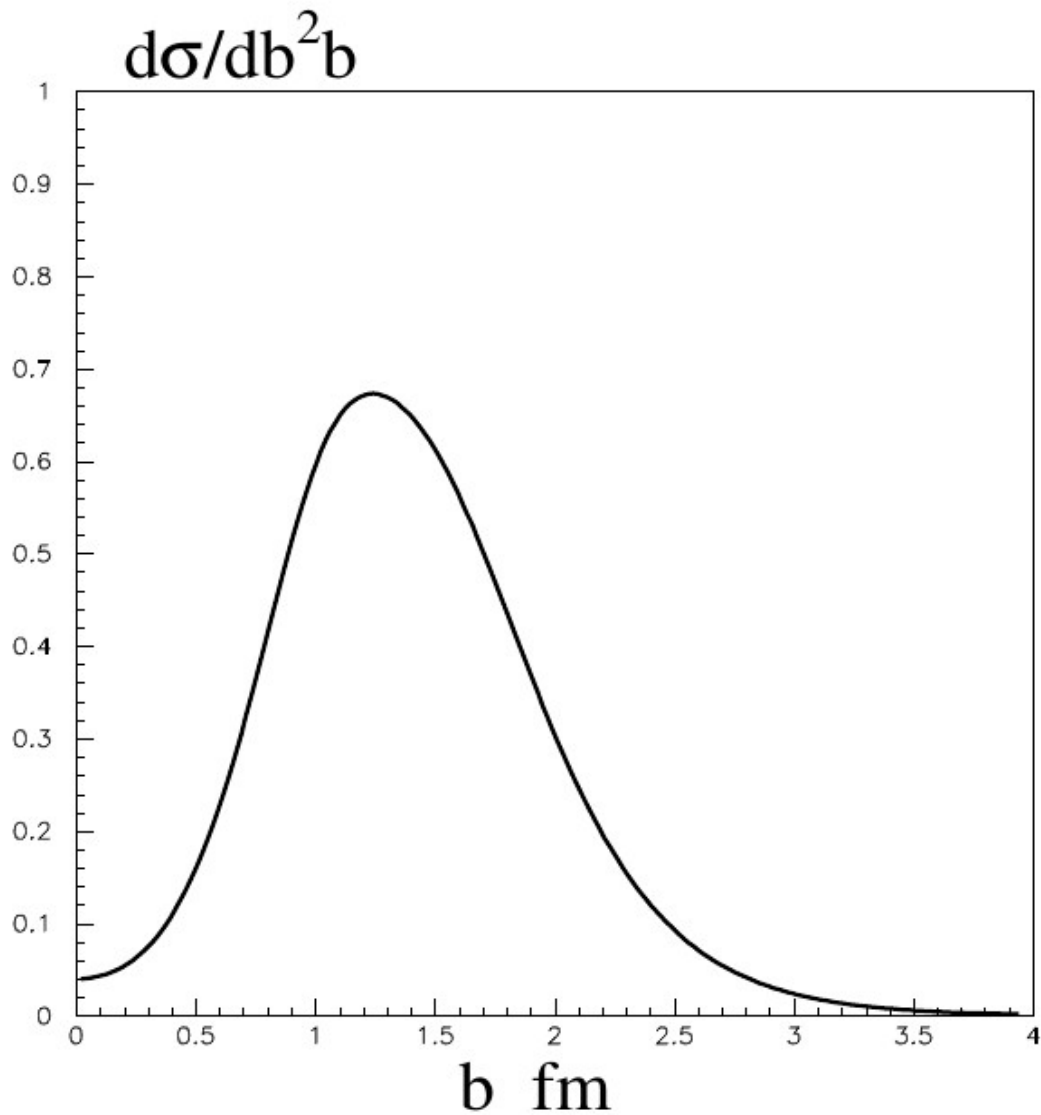
$$\alpha'_P = 0.052 \text{ GeV}^{-2} \quad g_N^2 = 23 \text{ mb}$$



Data - CMS 1503.08689



The model underestimates
CMS LRG cross sections
But overestimates the
ATLAS-ALFA SD



Due to small gap survival

$$S^2(b) = e^{-\Omega(b)}$$

at small b the LRG amplitude

becomes peripheral with

Maximum at $b \sim 1.2 - 1.3$ fm.

LRG in proton-lead collisions

a) incoherent

$$\sigma_{\text{incoh}} = \int d^2b T_e(b) |A|^2 S^2(b)$$

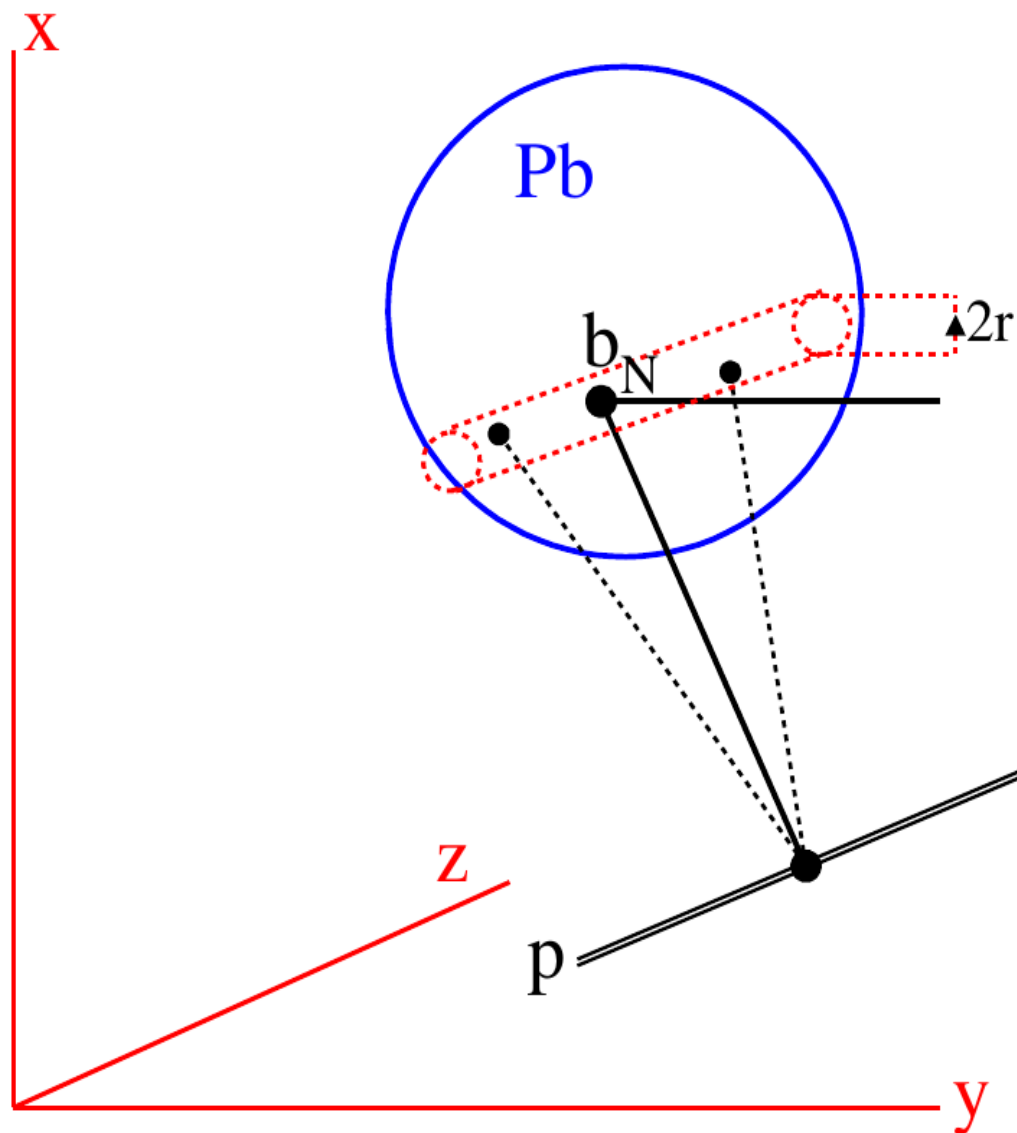
$$S^2(b) = \exp(-\langle \nu(b) \rangle)$$

$$\langle \nu(b) \rangle = T_s(b) \sigma(pN) - 1$$

$$T_e(b) = \frac{\int d^2b' T(b') |A(\vec{b} - \vec{b}')|^2}{\int d^2b' |A(b')|^2}$$

$$\sigma(pN) = \sigma_{\text{tot}} - \sigma_{\text{el}} - \sigma_{\text{LRG}}$$

$$\sigma(pN) = 65 \text{ mb}$$



$$T(b) = \int_{-\infty}^{+\infty} dz (\rho_p(r) + \rho_n(r)) \quad (11)$$

with $r = \sqrt{z^2 + r_t^2}$ and $r_t = b$. For the nucleon density in the lead $\rho(r)$ we use the Wood-Saxon form [\[12\]](#)

$$\rho_N(r) = \frac{\rho_0}{1 + \exp((r - R)/d)}, \quad (12)$$

where the parameters d and R respectively characterise the skin thickness and the radius of the nucleon density in the heavy ion; $r = (z, r_t)$. For ^{208}Pb we take the recent results of [\[13\]](#) [\[14\]](#)

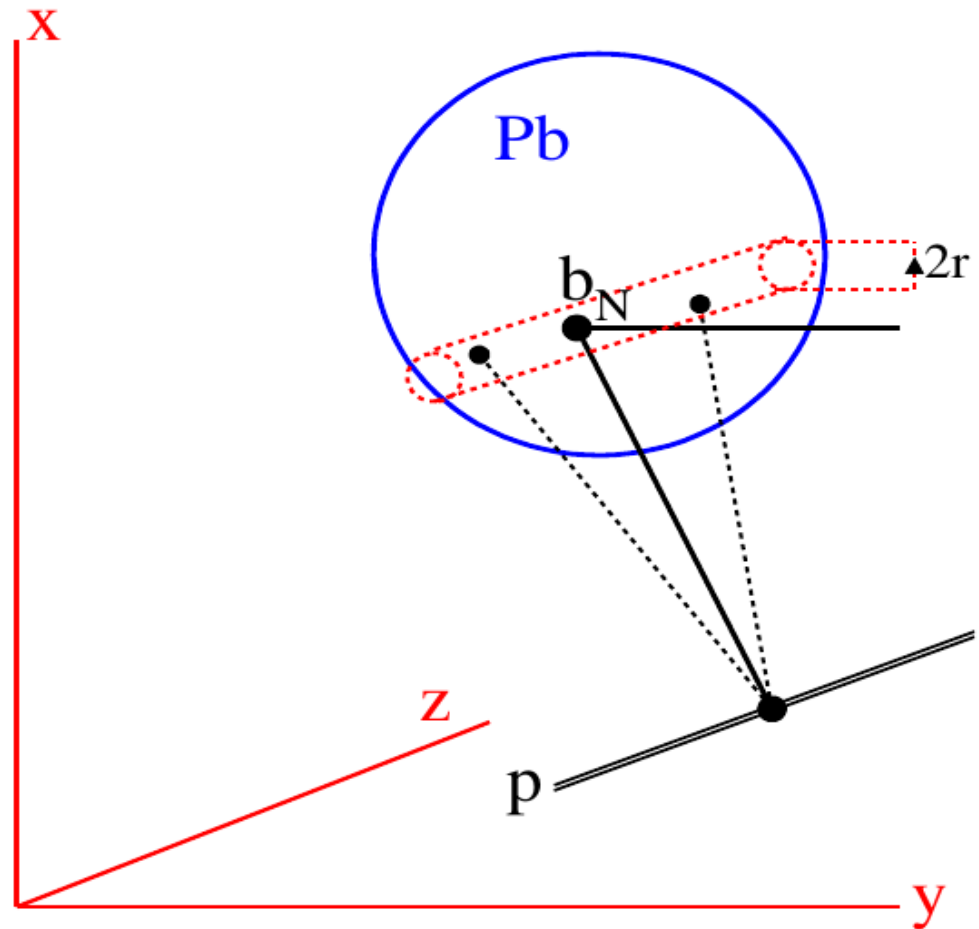
$$\begin{aligned} R_p &= 6.680 \text{ fm}, & d_p &= 0.447 \text{ fm}, \\ R_n &= (6.67 \pm 0.03) \text{ fm}, & d_n &= (0.55 \pm 0.01) \text{ fm}. \end{aligned} \quad (13)$$

$$T_e(b) = \frac{\int d^2b' T(b') |A(\vec{b} - \vec{b}')|^2}{\int d^2b' |A(b')|^2}$$

b) coherent

$$\sigma_{\text{coh}} = 4\pi B_{sd} F_{\text{Pb}}^2(t_{\text{min}}) \int d^2b T_e^2(b) |A|^2 S^2(b) ,$$

$$B_{sd} = 1 / \langle k_T^2 \rangle ,$$



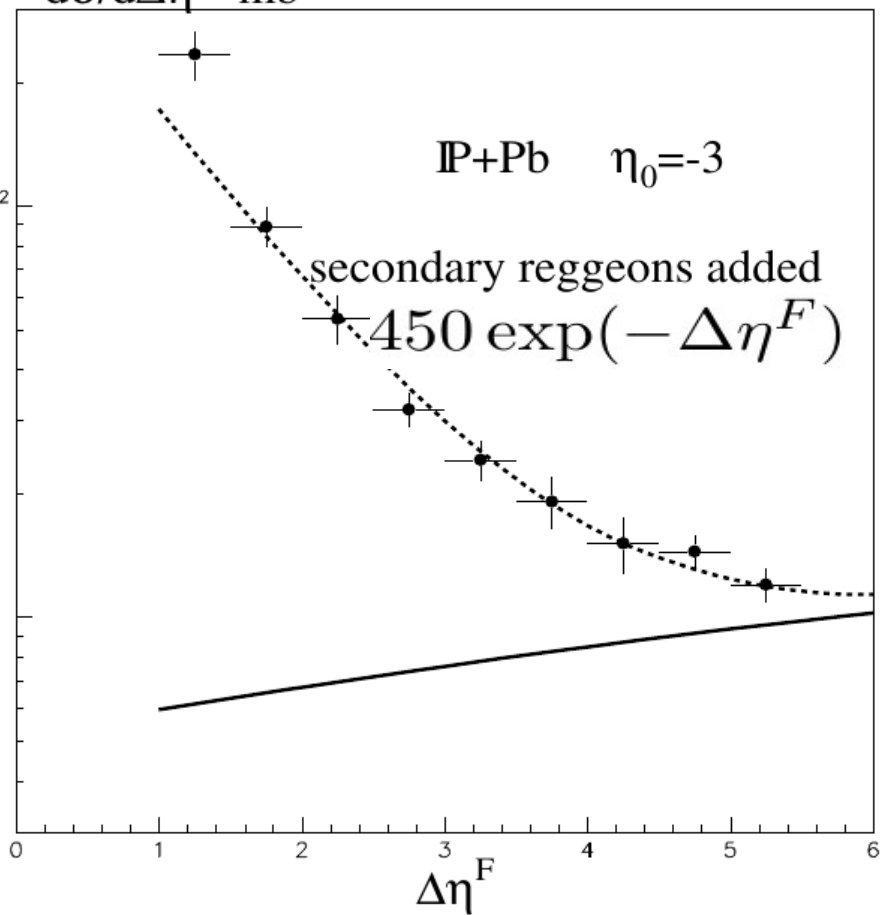
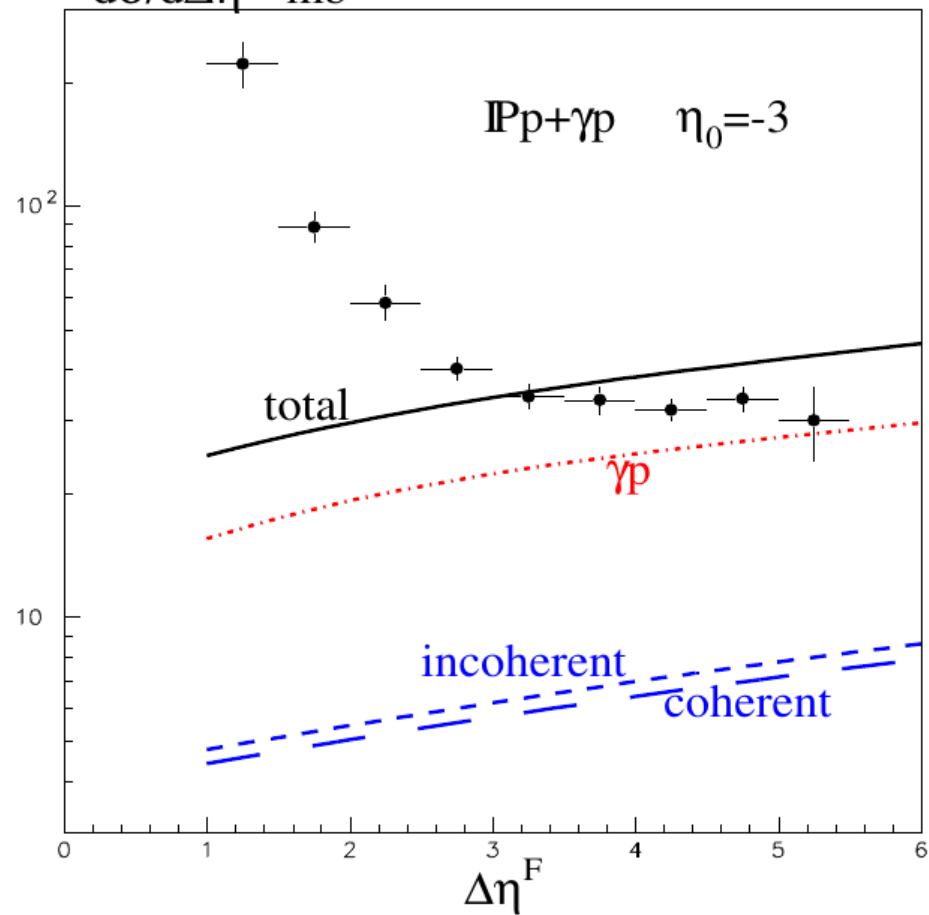
Photon exchange

$$\frac{x\sigma(Pb \rightarrow \gamma + p)}{dx} = \sigma(\gamma + p) \int d^2b \frac{xd^3n_\gamma(x, b)}{dx d^2b} \cdot S^2(b)$$

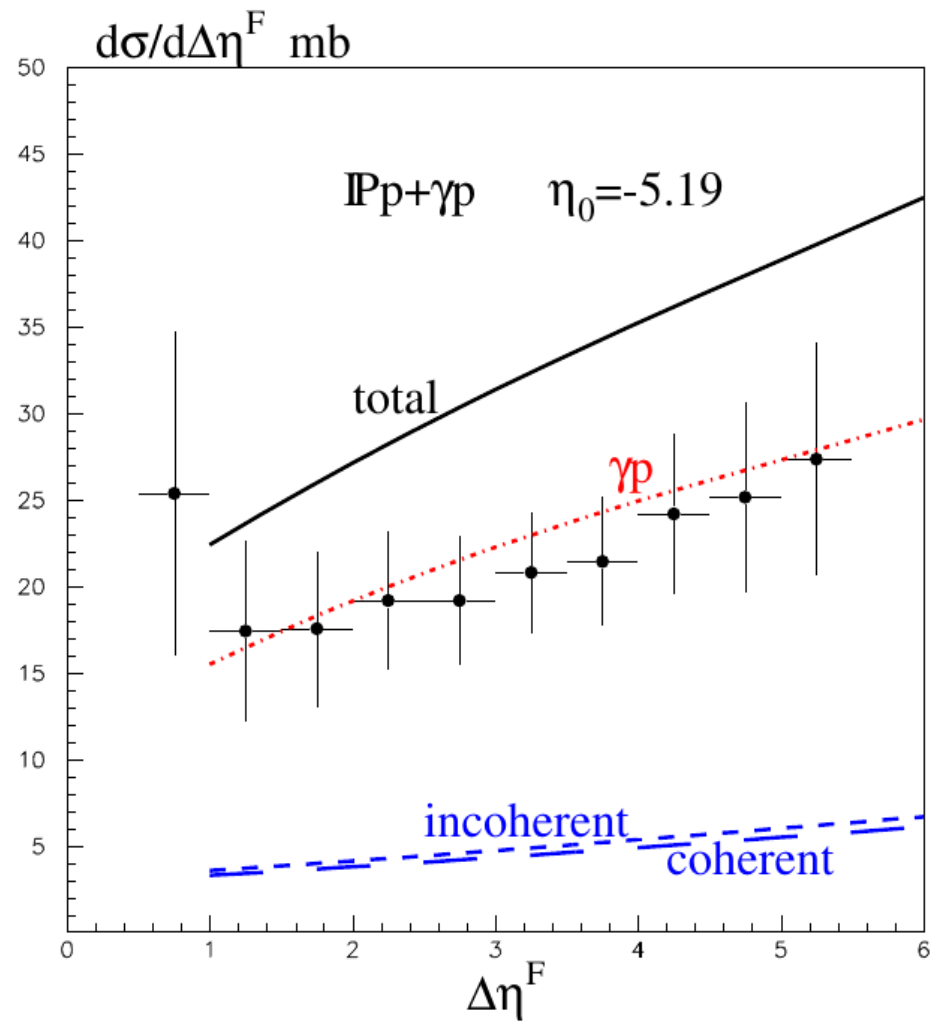
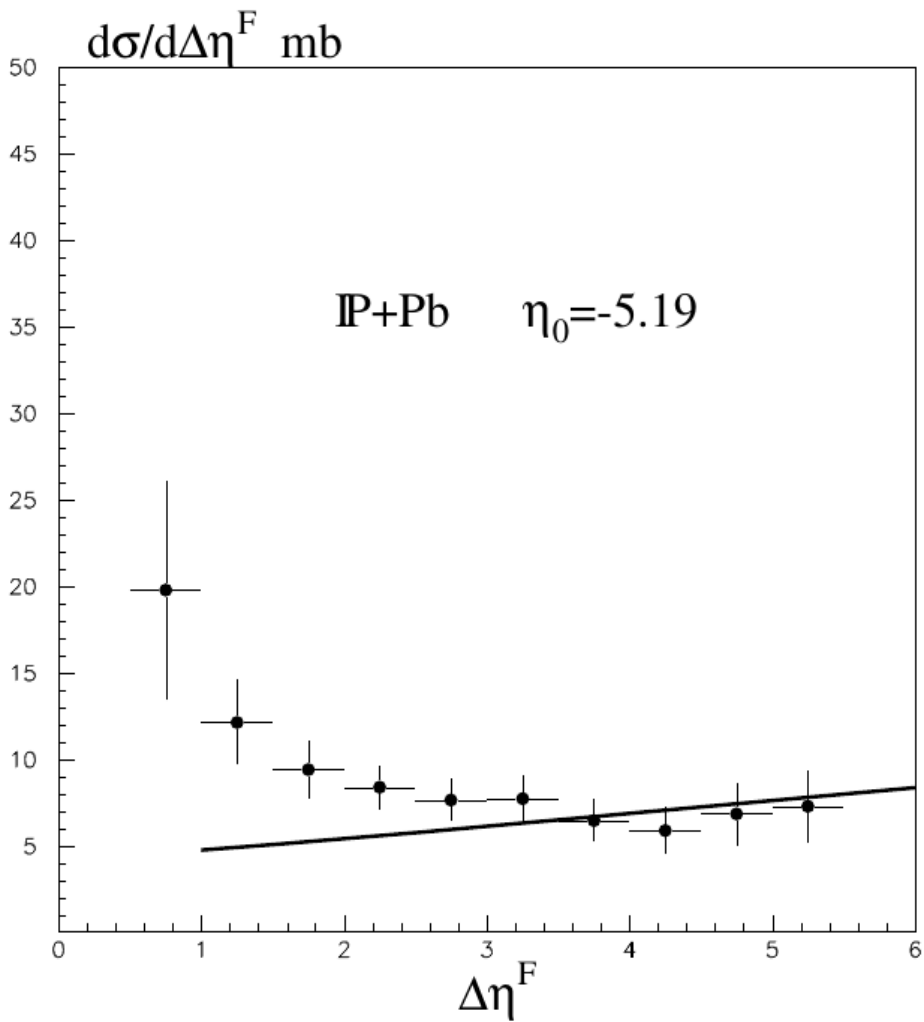
$$\frac{d^3n_\gamma}{dx d^2b_\gamma} = \frac{Z^2 \alpha^{\text{QED}}}{x\pi^2 b_\gamma^2} (xm_N b_\gamma)^2 K_1^2(xm_N b_\gamma) \quad Z^2 = 82^2 = 6724$$

$$\sigma(\gamma + p) = 0.0677 \text{mb} \cdot s_{\gamma p}^{0.0808} + 0.129 \text{mb} \cdot s_{\gamma p}^{-0.4525}$$

A. Donnachie, P.V. Landshoff, Phys.Lett. B **296** (1992) 227.

$d\sigma/d\Delta\eta^F$ mb $d\sigma/d\Delta\eta^F$ mb

No HF veto



Note that neglecting the transverse size of the elementary amplitudes both for the LRG amplitude and the absorptive cross-section we get almost twice the smaller cross-section – 4.6 mb instead of the 8.4 mb.

the result depends on the value of absorptive cross-section $\sigma(pN)$
replacing $\sigma(pN) = 65$ mb by 70 mb
changes the total value of $d\sigma/d\eta^F$ at $\Delta\eta^F = 6$
from 42.5 mb to 41.2 mb (right)
and from 8.4 mb to 7.8 mb (left)

If no additional experimental cuts were imposed to select the $\mathbb{P}p + \gamma p$ configuration, then we face the problem – why is the Pomeron-induced coherent plus incoherent contribution, which is expected to be about 13 mb, is not seen in the data?

THANK YOU