# **Rotation of Nuclear System**

# in Monte-Carlo Calculations of α-Particle and Fission Fragments Trajectories

## **Experimental Setup for Search of the TRI-Effect**





Search for a TRIple correlation B :  $B = (\sigma \cdot [p_{LF} \times p_{TP}])$ (note: all vectors are unit vectors)

Angular distribution of TPs :  $W(\theta) d\Omega \sim \{1 + D \cdot B(\theta)\} d\Omega$ where D measures size of correlation.

Experiment:  $D = (N_{+z} - N_{-z}) / (N_{+z} + N_{-z})$ 

Result : **count rates** for LF to the Left and TP upwards are **different** for  $s_z = +\frac{1}{2}\hbar$  and  $s_z = -\frac{1}{2}\hbar$ ,

### **TRI-effect**

Angular Distributions for Detector Combination: Left-Up. Red line -  $P_z>0$  (s<sub>z</sub>=+1/2ħ), blue line -  $P_z<0$  (s<sub>z</sub>=-1/2ħ), (D>0).





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# **Rotation motion in deformed nuclei**

$$E_{rot} = \frac{\mathbf{h}}{2\mathfrak{T}} \cdot (J(J+1) - K^2),$$
  
where J – total momentum,  
K – its projection  
\mathfrak{T} – moment of inertia

$$\mathbf{R} = \mathbf{J} - K \cdot \mathbf{n}$$
$$\mathbf{W}^2 \mathfrak{I}^2 = \mathbf{h}^2 (J(J+1) - K^2)$$

$$P(J^{+}) = \frac{2I+3}{3 \cdot (2I+1)} \cdot P_n \quad \text{for} \quad J^{+} = I + 1/2$$
$$P(J^{-}) = -\frac{1}{3} \cdot P_n \qquad \text{for} \quad J^{-} = I - 1/2$$

 $J^+ = I + 1/2$   $J^- = I - 1/2$ 

### Parameters for target nucleus <sup>235</sup>U:

$$----J = K + 2$$

-----J = K + 1

-----J = K

-----J = K + 3

The contributions of  $\sigma(J^+)$  and  $\sigma(J^-)$  to the fission cross-section:

Angular momentum of target nucleus I = 7/2

$$s(J^+ = 4) = 553 b$$
 and  $s(J^- = 3) = 323 b$   
 $P(J^+) = 5/12 \cdot P_n$   $P(J^-) = -1/3 \cdot P_n$ 

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The motion of the three fragments under the influence of their mutual Coulomb interaction cannot be calculated in closed form. The trajectories must therefore be calculated numerically.

The equations of motion are:

where

 $X_{ij}$  is the *j*th coordinate  $X_j$  of *i*th particle,  $V_{ij}$  the *j* component of the velocity  $V_i$  $F_{ij}$  the *j* component of the force  $F_i$  acting on the *i*th particle,  $m_i$  its mass.

These equations are replaced by the difference equations:

 $X_{ij}^{n+1} = X_{ij}^{n} + \widetilde{V}_{ij}^{n} \Delta t,$  $V_{ij}^{n+1} = \widetilde{V}_{ij}^{n} + \frac{1}{2m_i} F_{ij}^{n} \Delta t,$ 

 $\frac{dX_{ij}}{dt} = V_{ij}$ 

 $m_i \frac{dV_{ij}}{dt} = F_{ij}$ 

where  $\tilde{V}_{ij}^{n} = V_{ij}^{n} + \frac{1}{2m_i}F_{ij}^{n}\Delta t$ 

and  $F_{ij}^{n}$  is the *j* component of the force acting on particle *i* at the position  $X_{i}^{n}$ 

$$\mathbf{\hat{F}}_{i}^{n} = e^{2} Z_{i} \sum_{k=1}^{2} Z_{k} \frac{X_{i}^{n} - X_{k}^{n}}{\left| \frac{\mathbf{r}}{X_{i}^{n} - X_{k}^{n}} \right|^{3}}$$

The subscript k refers to the two other particles, and the superscript n refers to the value of the parameter after nth time interval.

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The size of time interval is not chosen to be constant. The total time  $t_n$  after *n* time intervals is an exponential function of *n*:

 $t_n = t_0 e^{na}$ ,

and hence the size of the *n*th time interval is given by

 $\Delta t_n = t_n - t_{n-1} = t_{n-1}(e^a - 1)$ 

The parameter  $t_0$  determines the accuracy of the calculation at the beginning of the trajectory (t = 0).

The parameter *a* determines the accuracy of the calculation at the end of the trajectory  $(t = \infty)$ .

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# The choice of initial parameters must led up to a coincidence experimental and calculated distributions:

- Light and heavy fragment's mass distributions
- I Total kinetic energy distribution
- $\alpha$ -particle energy distribution 6000.
- I Angular distribution of  $\alpha$ -particle





Angels of rotation of the Light fragment and α-particle<br/>in laboratory coordinates system.While system rotates α-particle carried along, but lags behind.







### Distance of ternary fission products from center of mass



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### **Results for <sup>235</sup>U**

I nfluence of linear dimensions of fission fragment's and  $\alpha$ -particle's detectors on angular distribution







# Thank you for your attention!